Numerical modeling and measurement techniques

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USPAS Course on collective effects

Thursday, 25.06.2009

One summary remark:

⇒ **Cures for coherent effects**

- ○ Impedance reduction (i.e. control and budget specification in the design phase, or identify and remove sources for running machines)
- ○ Since these effects are consequence of a resonant response to excitations on the beam natural frequencies, a spread in these frequencies in general helps
  - → use nonlinearities (e.g. sextupoles and octupoles) to increase the transverse detuning with amplitude against transverse instabilities
  - → use higher harmonic number rf-systems to enhance the spread in the synchrotron frequencies against longitudinal instabilities
- ○ Increase the longitudinal emittance (if possible), because the high density (in phase space) beams are more unstable
  - → this helps against both longitudinal and transverse instabilities
- ○ Use active feedback (also called damper)
  - ✓ system of pick-up + kicker that detects coherent motion and suppresses it
  - ✓ depending on the type of instability, it may be too demanding in terms of power or band-width. Easier against slow, low-frequency instabilities

⇒ **Two-stream phenomena are generally avoided by fighting the prime cause**

- ○ e.g., improve vacuum, use coated beam pipes with low secondary emission

Contents of this lecture:

⇒ **Numerical simulations for modeling of multi-particle effects**

- ○ the electromagnetic problem
    - → definition or calculation of the driving terms (field or particle distributions)
- ○ the beam dynamics problem
    - → put the driving terms previously calculated into the tracking of the beam particles and study the effects
    - → the simulation technique
- ○ some examples of simulations of single-bunch effects
    - ✓ head-tail instabilities
    - ✓ TMCI
    - ✓ longitudinal effects (use of the 2nd harmonic, potential well distortion, microwave instability)

⇒ **Examples of observations of coherent effects in existing accelerators and comparisons with simulations**

- ○ tune shift measurements
- ○ instabilities

How do we simulate numerically a multi-particle effect on a particle beam ? (1st step –**the electromagnetic problem**)

- **Space charge**:
  - relies on analytical formulae for ellipsoidal/Gaussian bunches
  - uses a Poisson solver to get the beam field
- **Impedance**. A reliable model for the ring impedance is needed
  - One part is the resistive wall component from the beam pipe (analytical)
  - The other part:

    ∗ It can be given as the sum of the individual contributions given by each accelerator component. These contributions, stored in databases, are previously calculated by means of

      ✓ electromagnetic codes for complex geometries, which can output the field maps of the given device when excited with a pulse
      ✓analytical formulae for simple geometries (e.g. tapers, steps)
      ✓bench measurements

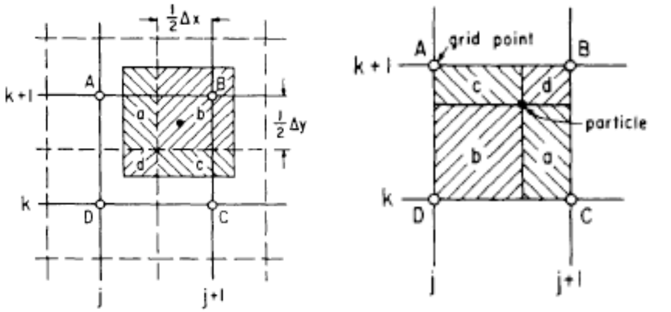
    ∗ It is the broad–band approximation of the accelerator
- **Two stream**:
  - relies on a numerical model of electron cloud formation/ion accumulation

How do we simulate numerically a multi-particle effect on a particle beam ? (2nd step –**the beam dynamics problem**)

- **Space charge**:
  - ✓ the additional space charge force is included in the single particle tracking by localizing it in some selected kick points along the lattice
- **Impedance**. Once the response of the ring to a pulse excitation is known, it can be used for calculating the corresponding kick on each particle of a bunch
  - ✓ single bunch effects have to be studied with full 6D bunches subdivided into longitudinal slices and calculating on each particle the effect of the kicks from the wakes of all preceding slices
  - ✓ multi bunch effects can be usually modeled with 4D bunches (x-y), which feel the effect of the wakes of all the preceding bunches
- **Two stream**:
  - ✓ electron cloud: beam particles are tracked through the accelerator and interact electromagnetically with an electron cloud lumped at some selected locations (single bunch)
  - ✓ ions: usually the ions are generated and tracked together with the beam particles (multi bunch)

The electromagnetic problem: **space charge**

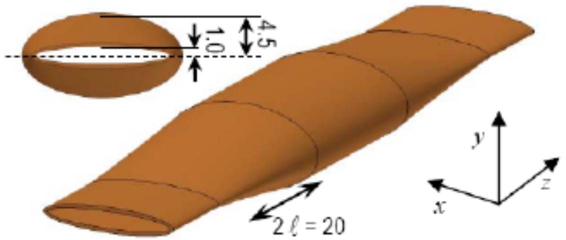
- The problem of the electromagnetic fields of some standard beam distributions in open space has been solved analytically for some cases. For example:
  - ✓ Ellipsoidal: R.W. Garnett and T.P. Wangler, 1981
  - ✓ Gaussian: M. Bassetti and G.A. Erskine. Closed expression for the electrical field of a two-dimensional Gaussian charge. CERN-ISRTH/80-06, 1980.
  - ✓ Formulae including the beam images for some standard chamber shapes, e.g. rectangular, also exist (see previous lecture)
- Poisson solvers for the general case
  - ✓ their input of the charge density is given by distributing the particles on a grid (usually with the Particle-In-Cell method)



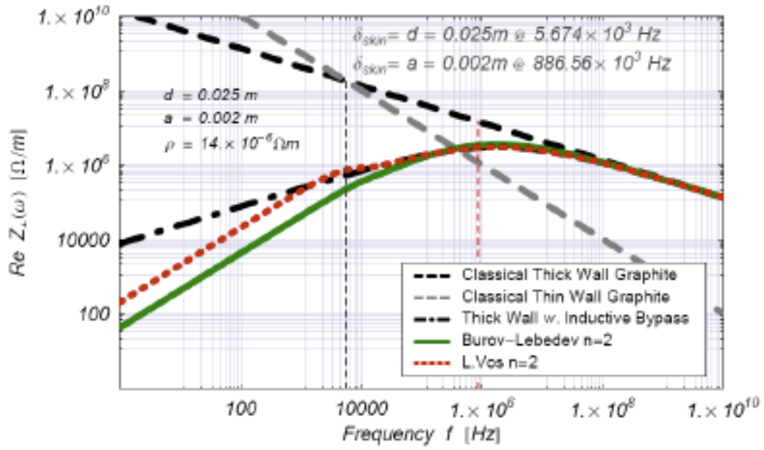
  - ✓ their solution includes the contribution of the images through the use of the appropriate boundary conditions
  - ✓ they can be based on solutions with the finite differences or FFT methods
  - ✓ they can have an adaptive grid and are usually very fast

The electromagnetic problem: **impedance (analytical)**

- Wake fields in relatively simple structures may be quite accurately obtained via analytical treatment leading to closed mathematical expressions.
- **Geometric effects** (induced by changes of cross-section, irises, cavities, etc., usually purely inductive impedances)
  - → Tapers in the inductive and diffractive regime, recently improved model w. r. t. the previous model by Yokoya and Stupakov
    - ✓ higher order terms included
    - ✓ elliptical cross-section
  - → Surface roughness
    - ✓ correlated and uncorrelated bumps
    - ✓ periodically corrugated structures



- **Resistive wall effects** (several regimes beyond the classical):
  - → long-range (low frequency, inductive by-pass)
  - → short-range (high frequency, ac conductivity)
  - → multi-layer boundary

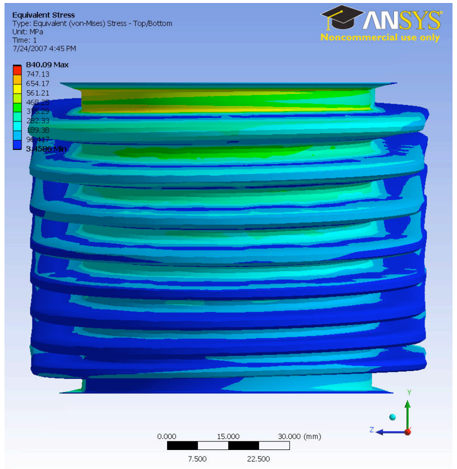


The electromagnetic problem: **impedance (numerical -1)**

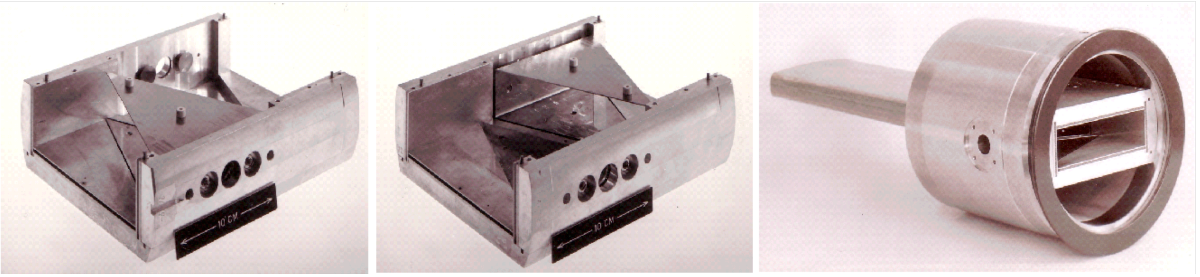
- Wake fields in a general structure may be most accurately obtained via numerical solution of Maxwell's equations.
- in the '80s the first **2D and 3D codes** were developed to solve numerically the Maxwell equations in given geometries (time or frequency domain)
  - → TBCI, MAFIA, ABCI, NOVO, XWAKE, ….
  - → More recently: GdfidL, HFSS, Microwave Studio, Particle Studio
- While newer rings built in the '90s tended to be based on a smooth design of the vacuum chamber such as to minimize geometric wakes from steps and abrupt transitions, they were made with flat/asymmetric chambers and shorter bunches (e.g. Linac based FELs):
  - → demand more powerful computation
  - → smaller mesh (often over a larger volume) & longer integration time
  - → larger memory and cpu time
- Many of these codes have been **parallelized** and can run on a cluster of cpu's
  - → GdfidL divides the integration space in sub-volumes, to be distributed over different nodes
  - → PBCI decomposes the computational volume with a load balancing scheme

The electromagnetic problem: **impedance (numerical -2)**

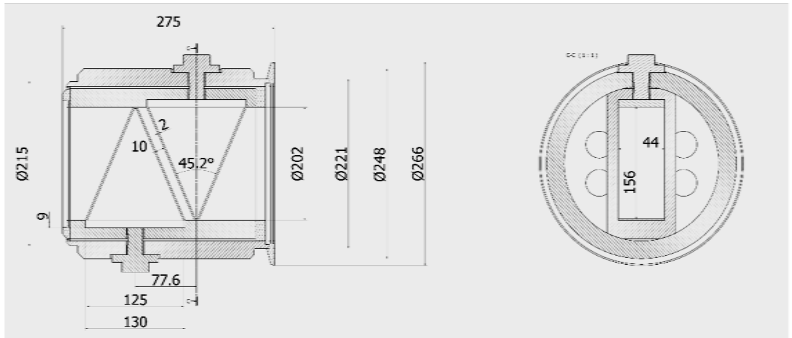
- Examples:
  - → Diagnostics equipments. For instance:
    - ✓ Wire scanners
    - ✓ Beam Position Monitors
  - → Kickers (injection, extraction, Q-measurement, dump)
  - → Collimators (betatron, energy), spoilers, scrapers
  - → Interconnectors, bellows



PS bellow

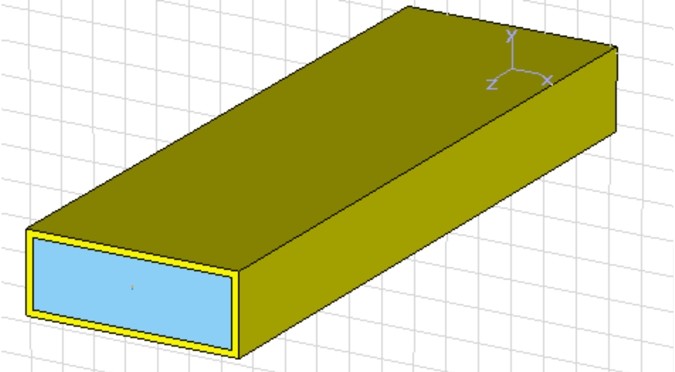


SPS BPMs



The electromagnetic problem: **impedance (numerical -3)**

- Example of use of Particle Studio:
  - → gives directly the wake field using a Gaussian bunch as source
  - → can be used for a simple structure for benchmark with theory



Geometric parameters

Thickness Copper = 0.2cm   1cm

Length = 1m 0.2m

Vacuum Chamber:

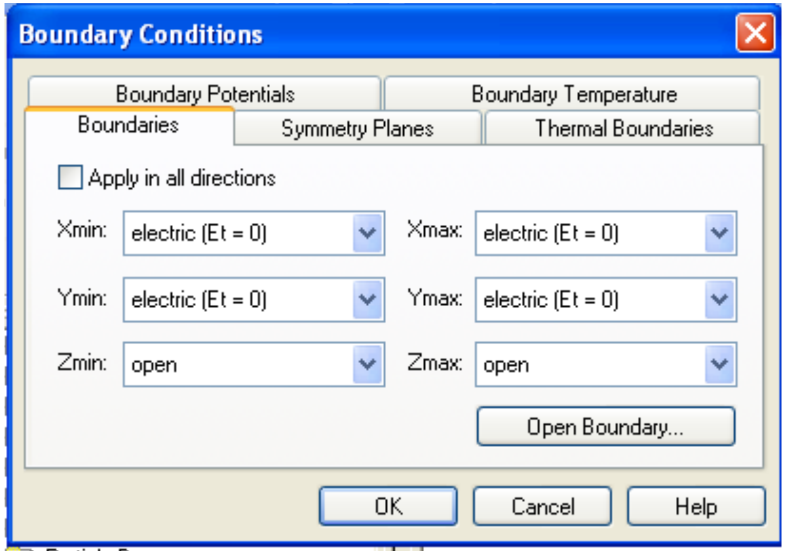
Rectangular shape : height=2cm; width= 6cm

Particle Beam Parameters

σ_{bunch} = 1cm, 0.8cm, 0.5cm

Charge = 1e-9

β=1

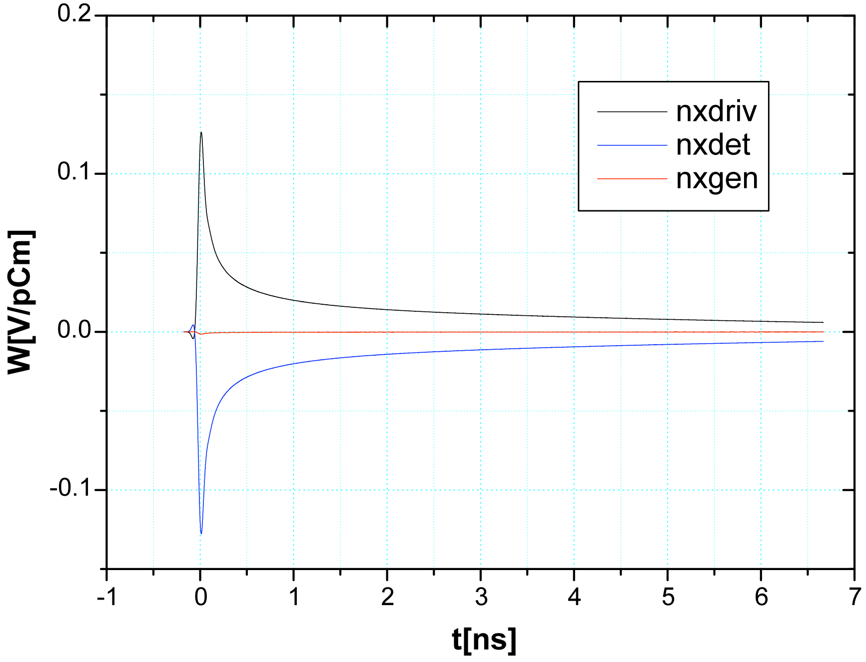


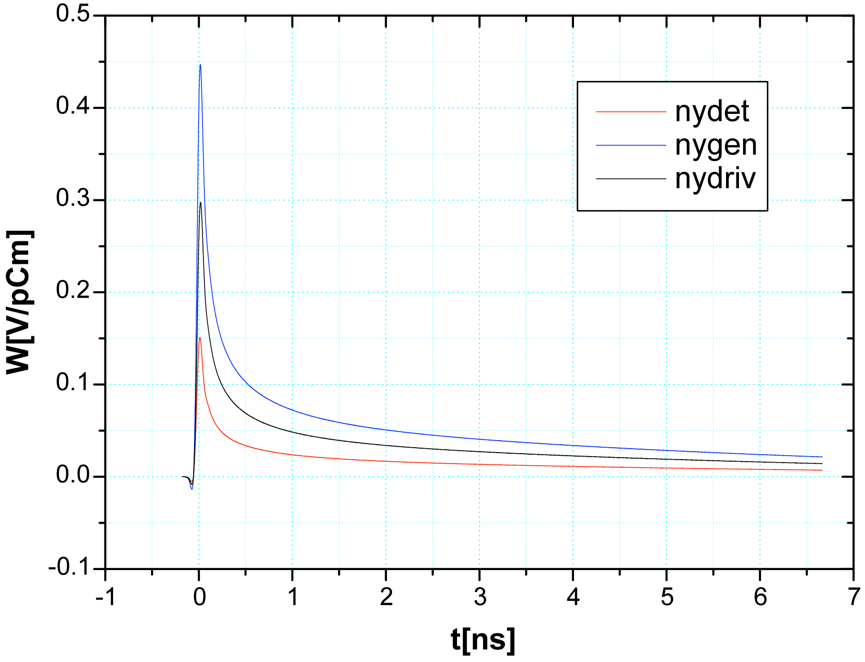
☐ Apply in all directions

Xmin: electric (Et = 0)
Xmax: electric (Et = 0)
Ymin: electric (Et = 0)
Ymax: electric (Et = 0)
Zmin: open
Zmax: open

The electromagnetic problem: impedance (numerical -4)

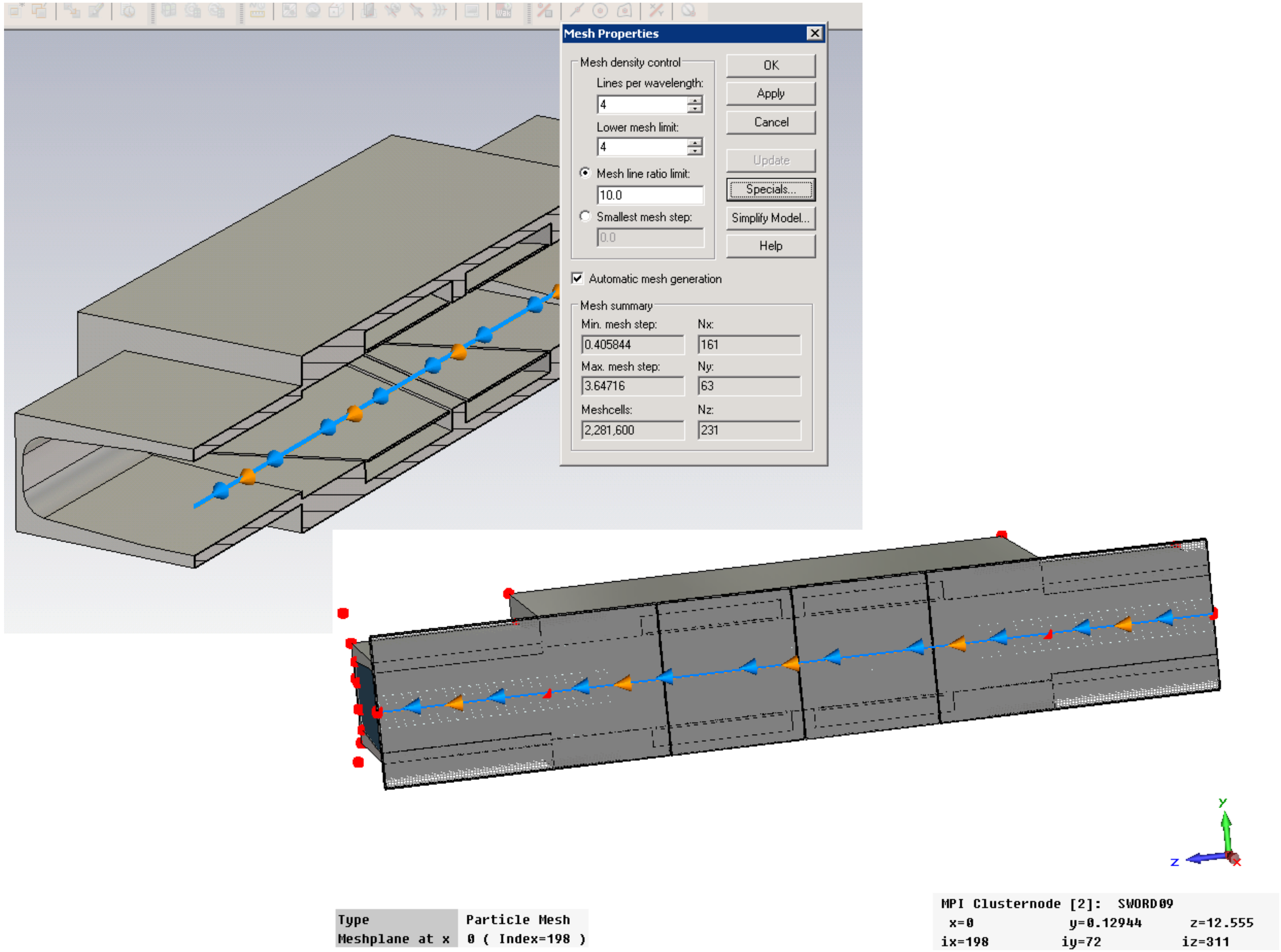
- Example of use of Particle Studio:
  - → with the previous structure we expect to see the resistive wall wake field
  - → since it is rectangular we could also disentangle dipolar and quadrupolar wakes
  - → as expected from the chosen aspect ratio, the Yokoya coefficients are recovered





The electromagnetic problem: **impedance (numerical -5)**

- Example of use of Particle Studio:
  - → More complicated structures can be simulated, e.g. the SPS-BPMs



Mesh density control

Lines per wavelength: 4

Lower mesh limit: 4

● Mesh line ratio limit: 10.0

○ Smallest mesh step: 0.0

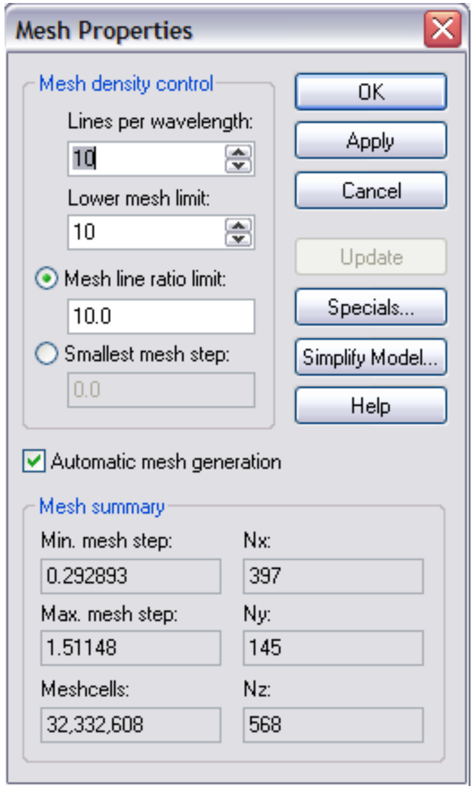
☑ Automatic mesh generation

Mesh summary

Min. mesh step: 0.405844	Nx: 161
Max. mesh step: 3.64716	Ny: 63
Meshcells: 2,281,600	Nz: 231

```
Type            Particle Mesh
Meshplane at x  0 ( Index=198 )
```

```
MPI Clusternode [2]:  SWORD09
 x=0            y=0.12944       z=12.555
ix=198          iy=72           iz=311
```



Mesh density control

Lines per wavelength: 10

Lower mesh limit: 10

● Mesh line ratio limit: 10.0

○ Smallest mesh step: 0.0

☑ Automatic mesh generation

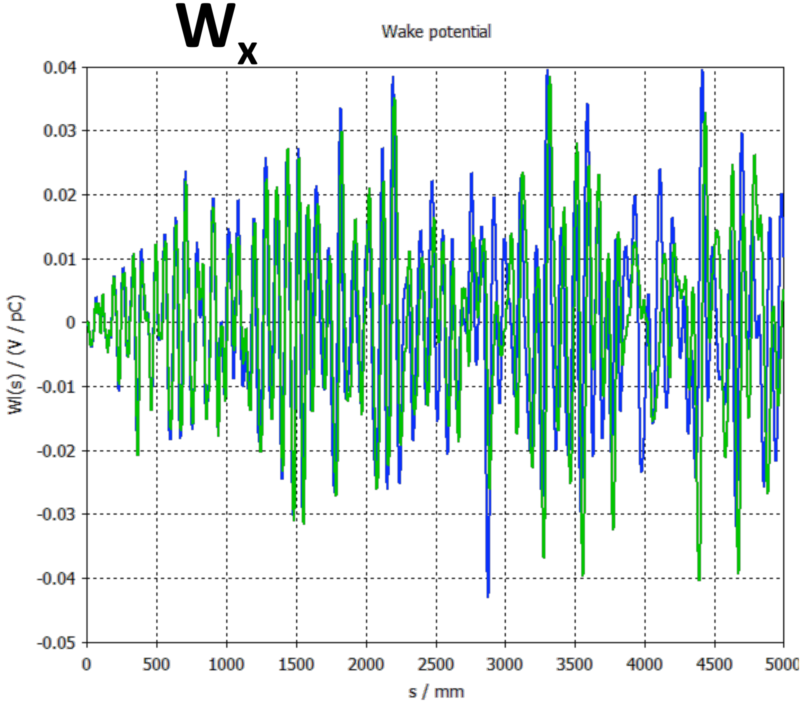
Mesh summary

Min. mesh step: 0.292893	Nx: 397
Max. mesh step: 1.51148	Ny: 145
Meshcells: 32,332,608	Nz: 568

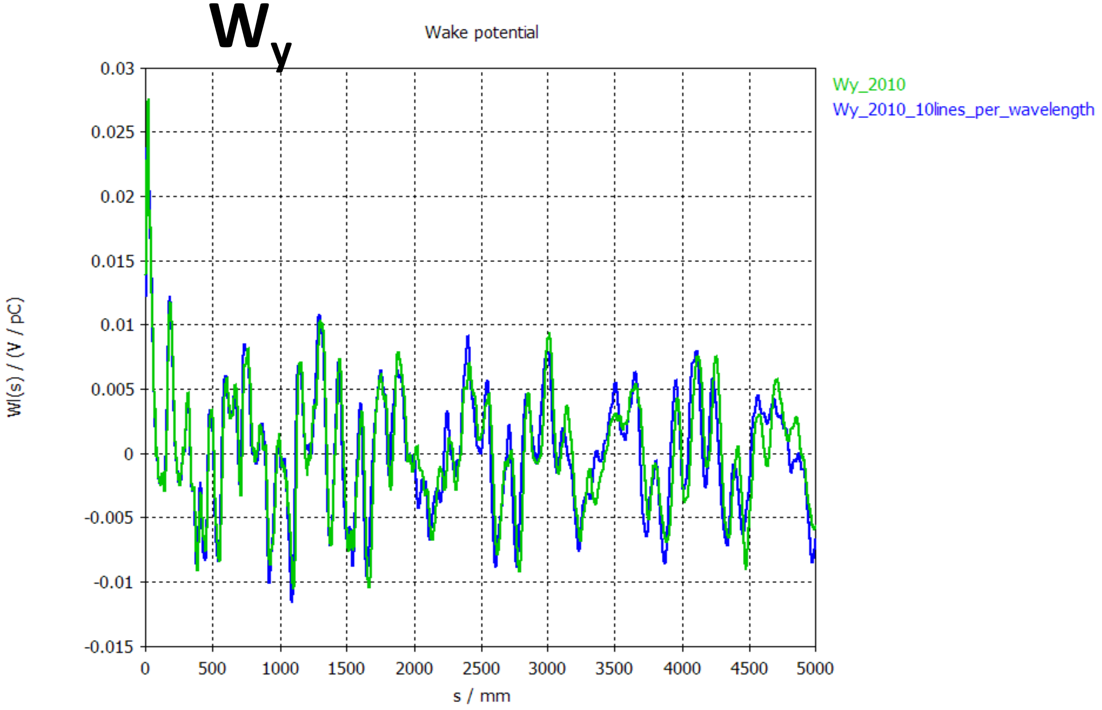
The electromagnetic problem: **impedance (numerical -6)**

- Example of use of Particle Studio:
  - → More complicated structures can be simulated, e.g. the SPS-BPMs

W_x



W_y



MovieEx, MovieEy, MovieEz, MovieEz2

The electromagnetic problem: **impedance (bench)**

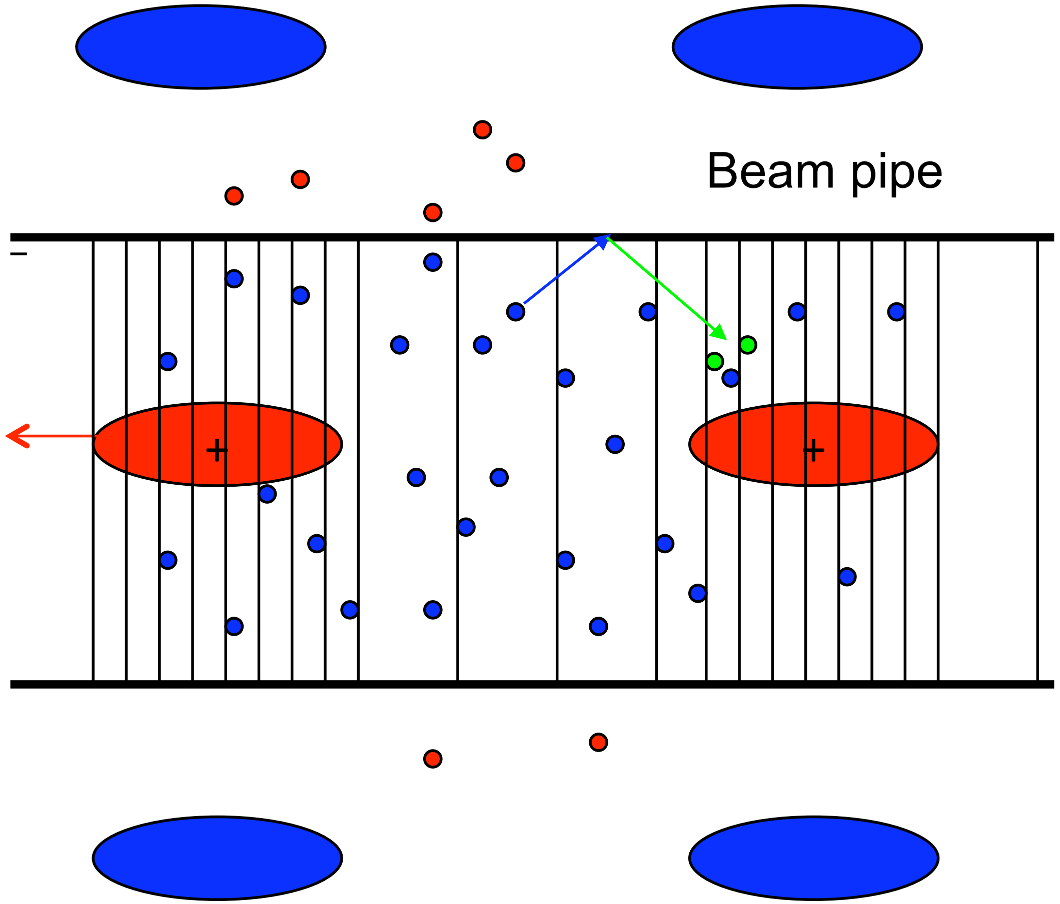
- Some devices can be tested in lab and their impedance is estimated from the scattering coefficients obtained with the 1- or 2- wire method. For example:
  - → Tubes (shielded, coated, grooved)
  - → Collimators (betatron, energy)
  - → Kickers

RF shielded ceramic pipe for RCS
Courtesy YH. Chin, J-PARC

LHC collimator prototypes in copper and graphite

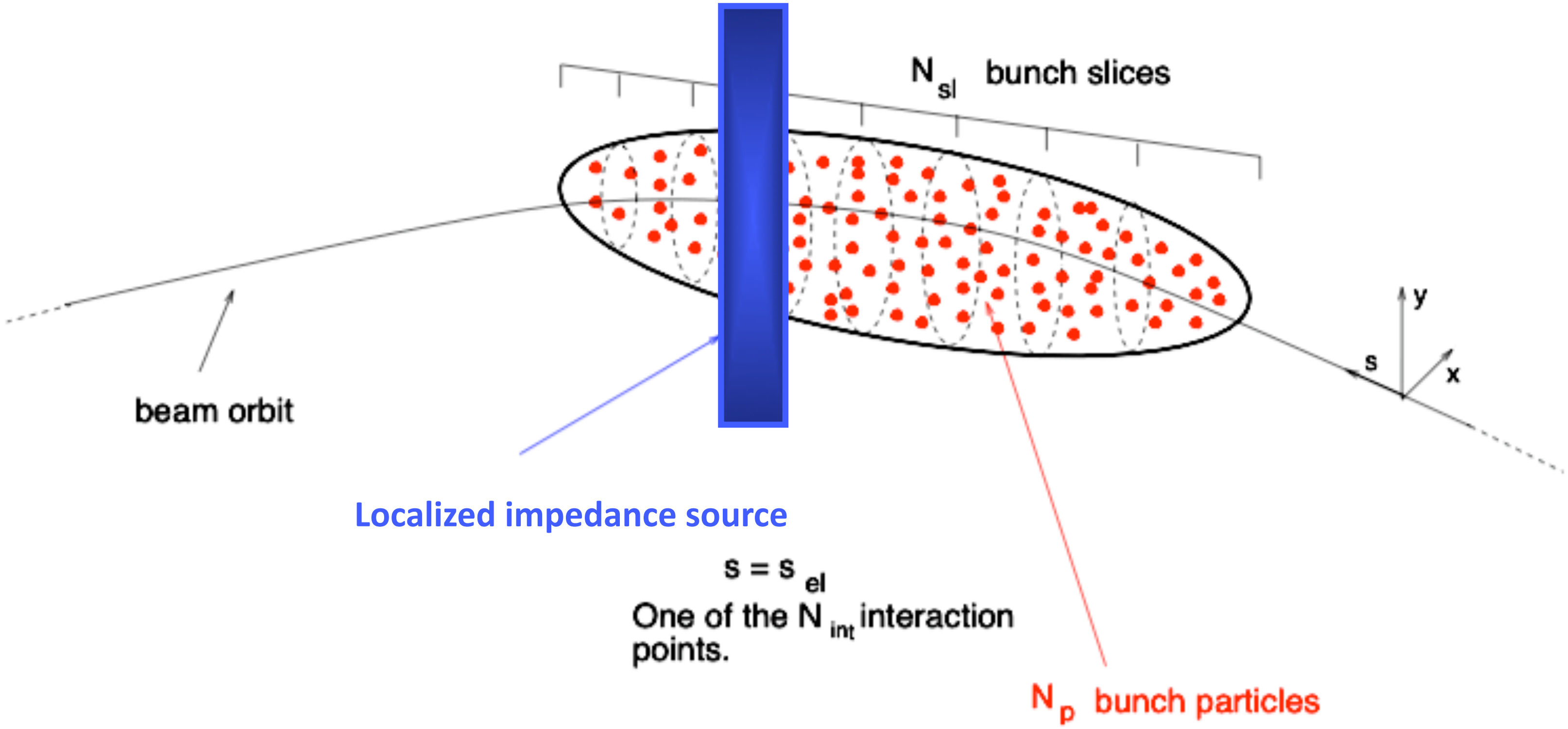
The electromagnetic problem: **two-stream (electron cloud)**

- To study the effect on the beam, we first need to model the electron cloud formation (ECLOUD code, F. Zimmermann et al.)



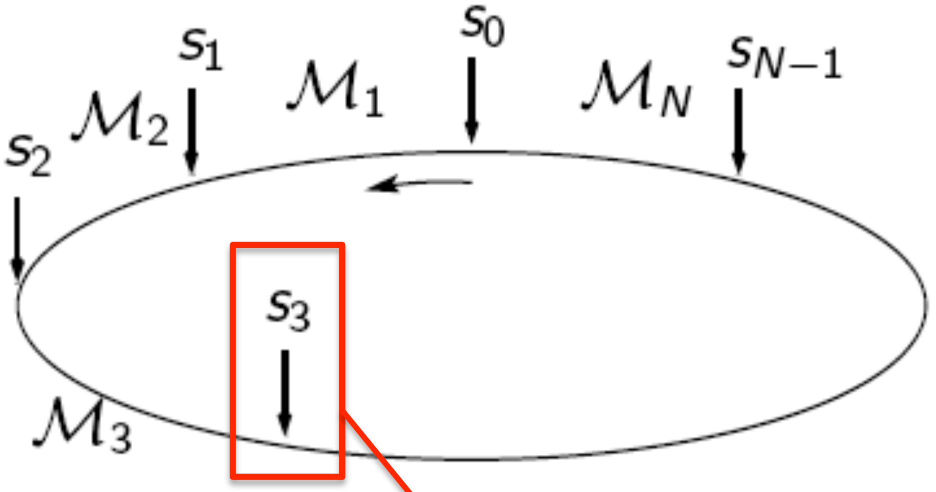
- focus on a beam line section (1m for ex.)
- slice bunch and interbunch gaps
- Electrons are macroparticles: they are created (photoemission or gas ionization) and accelerated in beam and image fields
- if the e- hits the wall create secondaries by changing its charge.
- After many bunches, the electrons come to a dynamic „steady“ state

The beam dynamics problem: **The physical model for single bunch (HEADTAIL)**



The collective interaction is lumped in one or more points along the ring (**kick points**), where the subsequent slices of a bunch (macroparticles) interact with an impedance (through the wake) or with an electron cloud

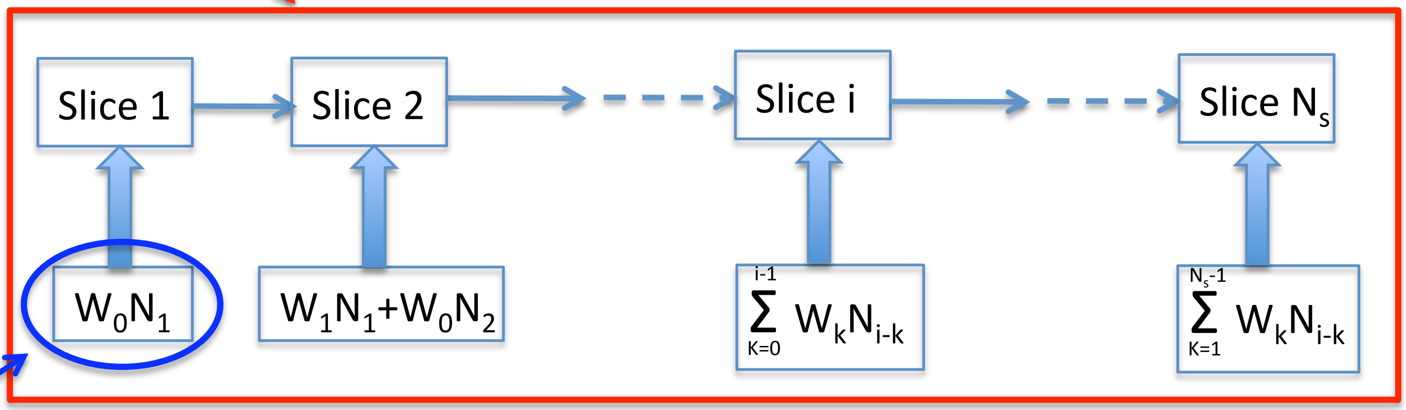
The beam dynamics problem: Numerical implementation (wake fields)



1. Bunch macroparticles are transported across different interaction points through the sector matrices
2. At each interaction point macroparticles in each slice receive the kick from the wakes of the preceding slices
3. Slicing is refreshed at each turn taking into account the longitudinal motion

Longitudinal

$W_i = W_{||}(i\ \Delta z)$



Slice 1 → Slice 2 → - - → Slice i → - - → Slice N_s

W_0N_1

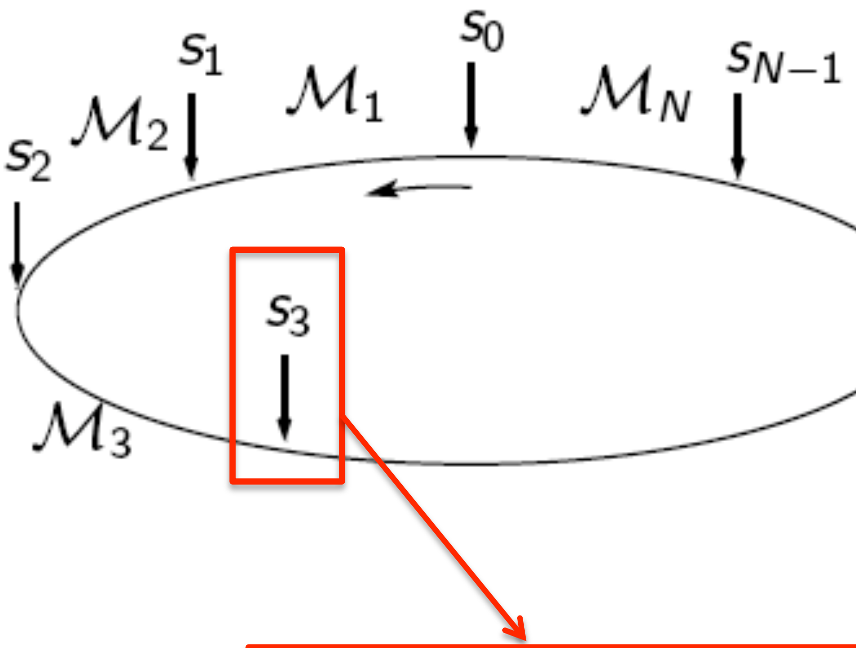
$W_1N_1+W_0N_2$

$\sum_{K=0}^{i-1} W_kN_{i-k}$

$\sum_{K=1}^{N_s-1} W_kN_{i-k}$

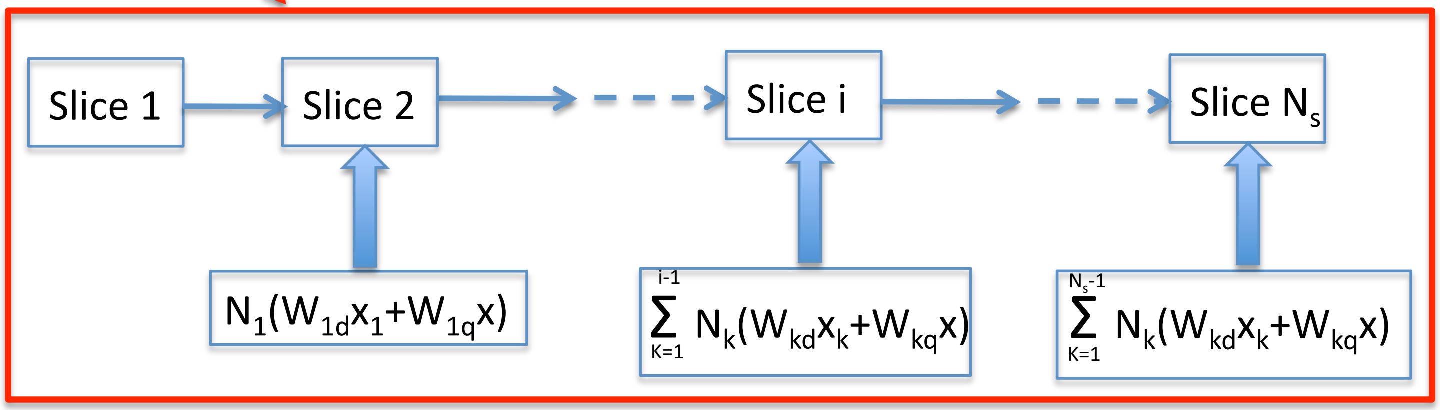
Energy loss

The beam dynamics problem: **Numerical implementation (wake fields)**



1. Bunch macroparticles are transported across different interaction points through the sector matrices
2. At each interaction point macroparticles in each slice receive the kick from the wakes of the preceding slices
3. Slicing is refreshed at each turn taking into account the longitudinal motion

Transverse (x)
dipolar:
$W_{id} = W_{dx}(i\,\Delta z)$
quadrupolar:
$W_{iq} = W_{qx}(i\,\Delta z)$
x_i centroid of slice i
x position of particle



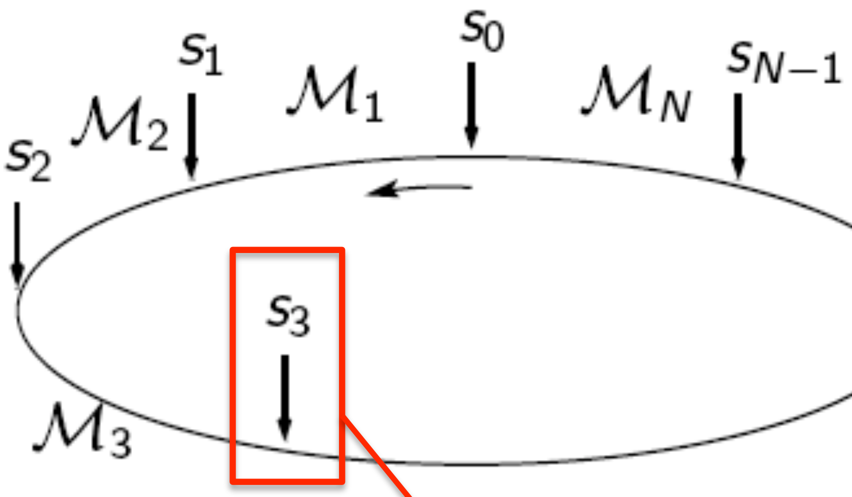
Slice 1 → Slice 2 → - - → Slice i → - - → Slice N_s

$N_1(W_{1d}x_1+W_{1q}x)$

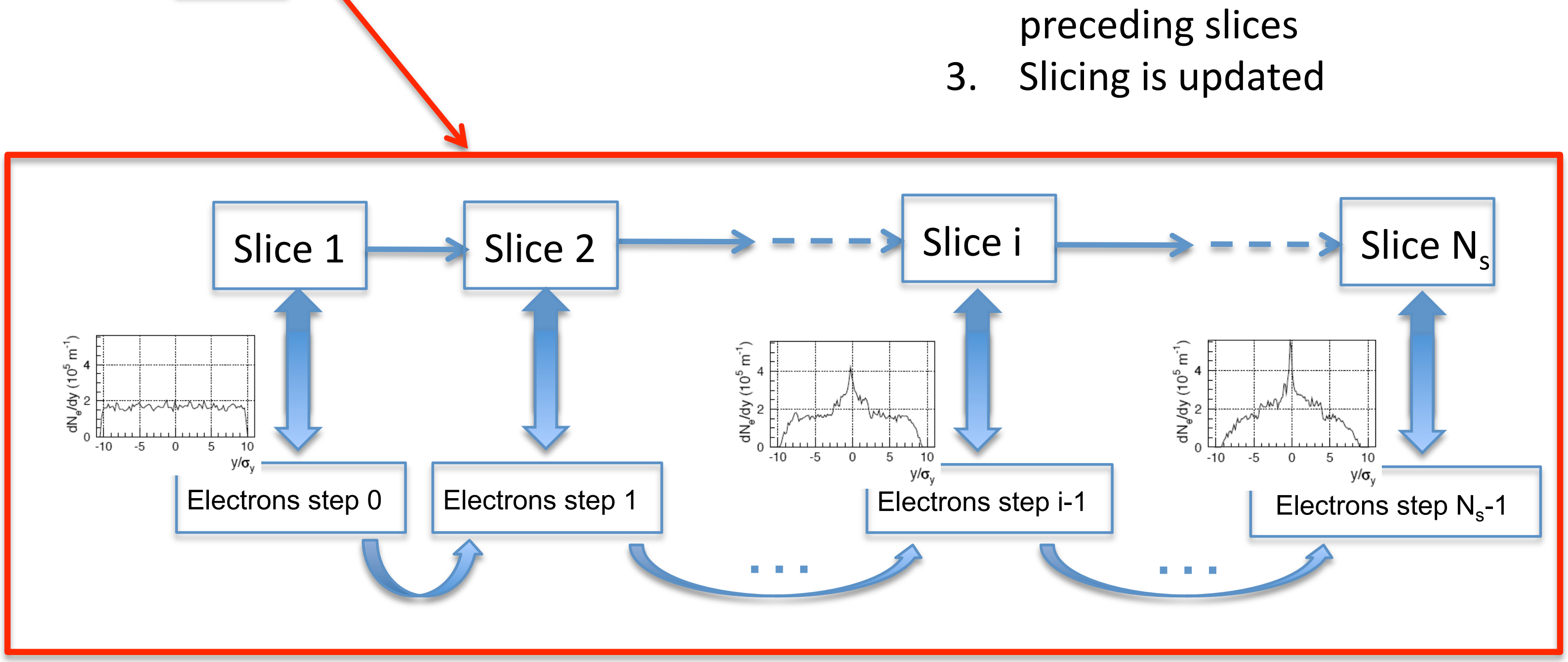
$\sum_{K=1}^{i-1} N_k(W_{kd}x_k+W_{kq}x)$

$\sum_{K=1}^{N_s-1} N_k(W_{kd}x_k+W_{kq}x)$

The beam dynamics problem: **Numerical implementation (electron cloud)**



1. Bunch macroparticles are transported across different interaction points through the sector matrices
2. At each interaction point macroparticles in each slice interact with the electron cloud, as it was modified by the interaction with the preceding slices
3. Slicing is updated



Slice 1

Slice 2

Slice i

Slice N_s

Electrons step 0

Electrons step 1

Electrons step i-1

Electrons step N_s-1

Features included in the HEADTAIL model (I)

- Full transverse and longitudinal motion
- Transverse motion modeled through a single turn matrix or transporting each particle's coordinates from one interaction point to the next one by using the correct transport matrices from MAD-X
- Synchrotron motion can be
  - Linear
  - Sinusoidal voltage
  - With an optional second rf system that can be switched on during the simulation
  - Inside an accelerating bucket, with or without higher order terms of η
  - Debunching (rf off)
  - Periodic over the circumference (coasting beam)
- Bunch initial distribution can be
  - Longitudinally: Gaussian or uniform
  - Transversely: Gaussian
- Chromaticity in both planes
- Detuning with amplitude
- Linear coupling

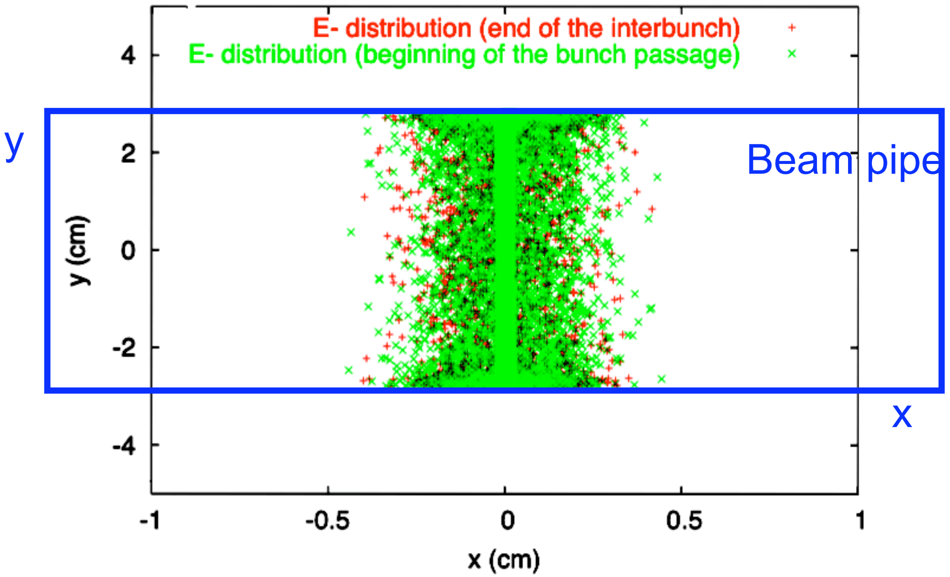
Features included in the HEADTAIL model (II)

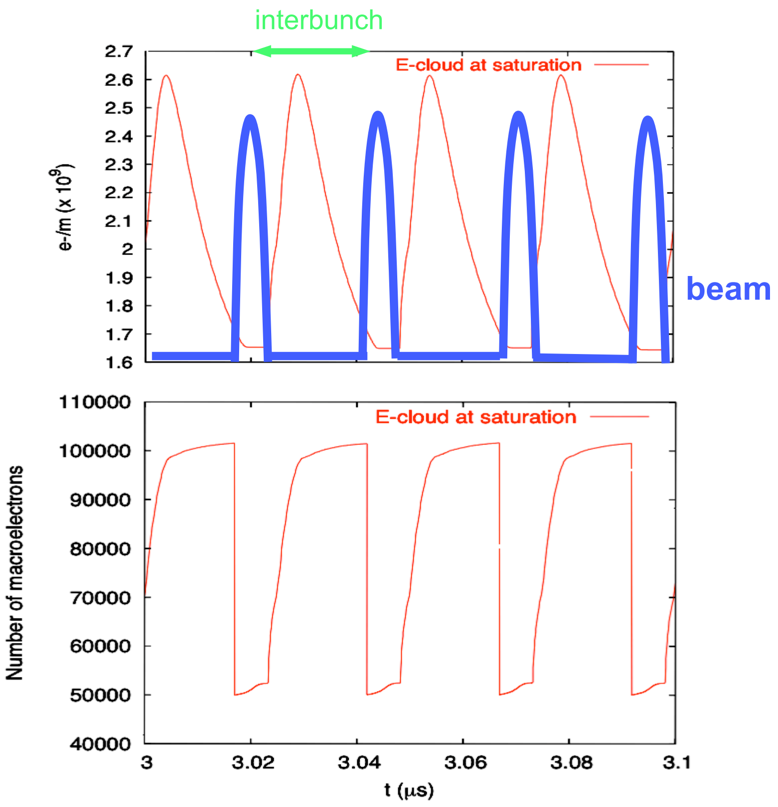
- Electron cloud kick(s):
  - Soft Gaussian approach with finite size electrons (used till 2002, obsolete)
  - PIC module on a grid inside the beam pipe
  - PIC solver with optional conducting boundary conditions
  - Uniform or 1-2 stripes initial e-distributions
  - Kicks can be given at locations with different beta functions and different electron cloud initial distributions (densities)
  - Electrons can move in
    - field free space
    - dipole
    - solenoid
    - combined function magnet
  - The initial distribution of electrons can be optionally loaded from the output of the ECLOUD code, which can save the exact electron distribution at saturation, right before the bunch passage

Quasi-selfconsistent model of electron cloud

→ Electron distribution used in HEADTAIL generally was uniform in the beam pipe

→ Model can be improved by using as an input the distribution of electrons at the beginning of a bunch passage, as it comes out of the build up ECLOUD code





Features included in the HEADTAIL model (III)

- Short range wake field is from
  - a broad band impedance

$$Z_{1\perp}(\omega) = \frac{\omega_R}{\omega} \frac{Z_\perp}{1 + iQ_\perp\left(\frac{\omega_R}{\omega} - \frac{\omega}{\omega_R}\right)}$$

  - Classical thick resistive wall.
  - Resistive wall with inductive by-pass
  - Loaded from external table
- Space charge. Optionally, each bunch particle can receive:
  - a transverse kick proportional to the local bunch density around the local centroid
  - a longitudinal kick proportional to the local derivative of the beam line density

$$\begin{pmatrix} x_{n+1} \\ x'_{n+1} \end{pmatrix} = M_1(\delta p) M_2(I_x, I_y) \left[M_{sc}(z) \begin{pmatrix} x_n - \hat{x}(z) \\ x'_n + \Delta x'_{EC, Z_{1\perp}} - \hat{x}'(z) \end{pmatrix} + \begin{pmatrix} \hat{x}(z) \\ \hat{x}'(z) \end{pmatrix} \right]$$

Example of simulation/measurements: **the head-tail instability**

- Due to chromaticity, single bunches develop head-tail modes ($m=1$), which can be strongly unstable at high intensity. The most dangerous mode is the mode $l=0$:

  → It is unstable below transition ($\gamma < \gamma_t$), if the chromaticity is positive ($\xi_{x,y} > 0$)

  → It is unstable above transition ($\gamma > \gamma_t$), if the chromaticity is negative ($\xi_{x,y} < 0$)

- Higher order modes ($l\geq1$) are unstable for negative chromaticities below transition and for positive chromaticities above transition. However, they are much slower and they can be naturally damped by other sources of tune spread, or can be suppressed with a damper.

- As a consequence, it is critical to control the mode $l=0$ by operating the machine with the correct sign of chromaticity.

  → Machines that run always below their transition energy (usually hadron machines) must have negative chromaticity (e.g., the CERN-PSB, GSI-SIS) and they can live with their natural chromaticity, which is negative for a classical lattice design. These machines can also avoid to use sextupoles for chromaticity correction

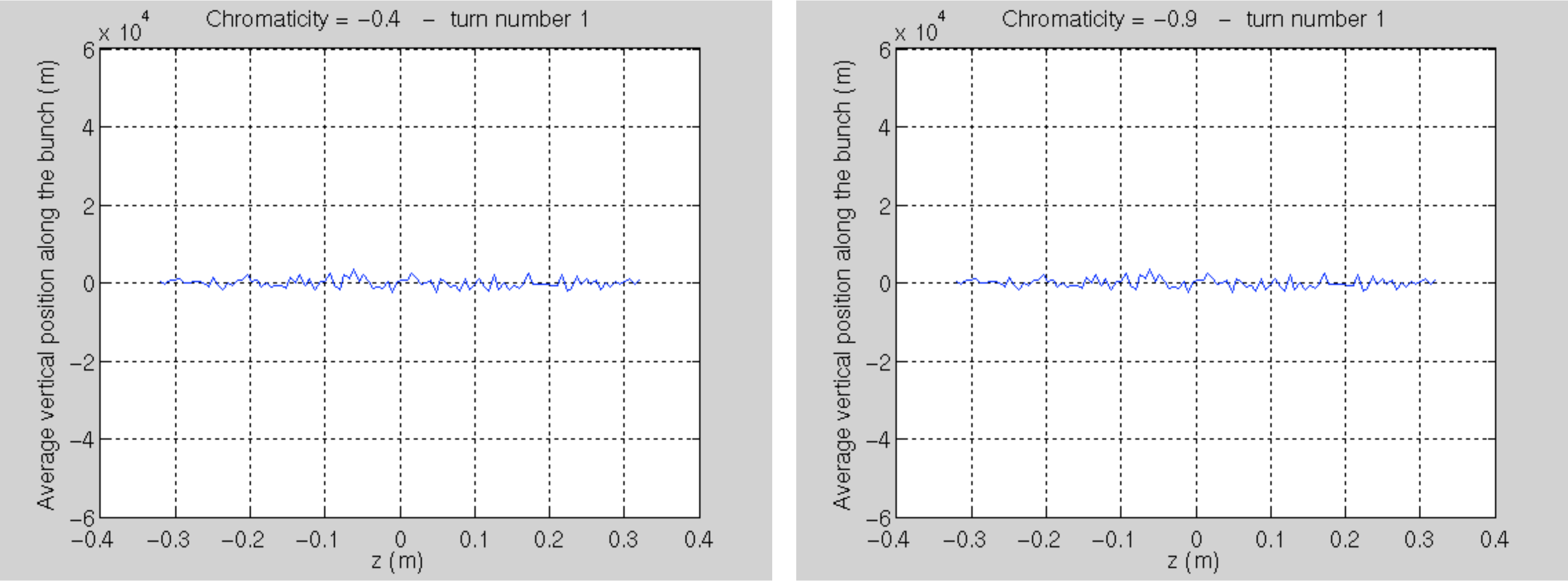
  → Machines that run always above transition energy (lepton machines, CERN-LHC, BNL-RHIC with protons) need chromaticity correction (and therefore two families of sextupoles) in order to make their chromaticity slightly positive.

  → Machines that cross transition (CERN-PS, CERN-SPS, BNL-RHIC with ions) need a scheme of synchronized swap of the sign of chromaticity at transition crossing

Example of simulation: **the head-tail instability**

⇒ The fundamental mode of a head-tail instability (m=1, *l=0*) can be simulated to have a detailed look at the instability evolution for different chromaticity values (assuming the SPS parameters and a simple broad band model for the impedance)

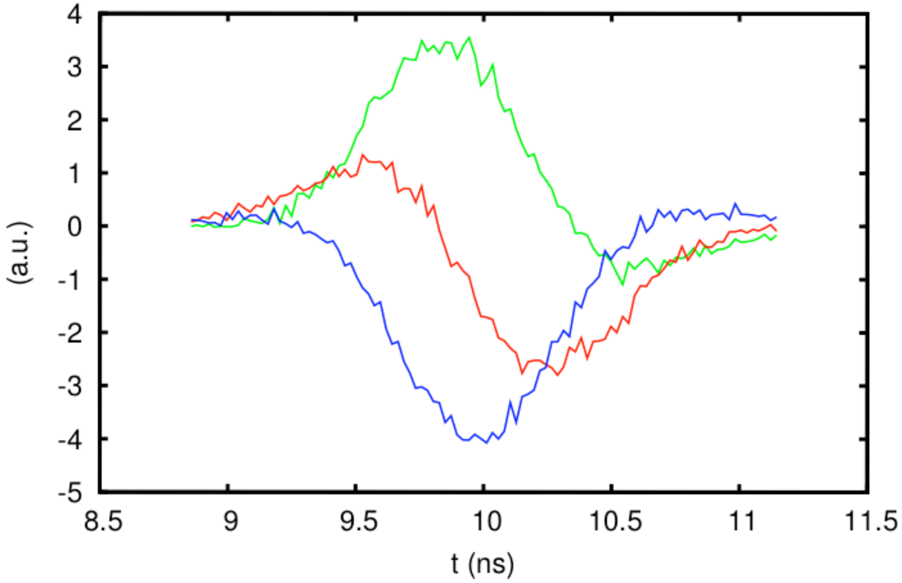
⇒ Movies show the evolution of the Δ (centroid) signal along the bunch over 1045 turns of unstable evolution for two chromaticity values (-0.4 and -0.9)



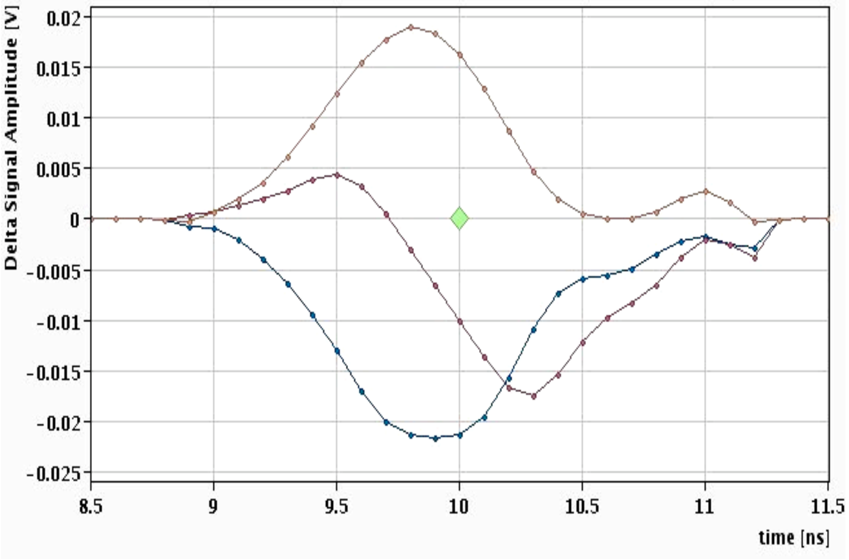
Example of simulation/measurements: **the head-tail instability**

- The fundamental mode of a head-tail instability can be simulated to have a detailed look at the instability evolution for different chromaticity values (assuming the SPS parameters and a simple broad band model for the impedance)
  - ⇒ The comparison between measurement and theory is impressive!
  - ⇒ Plots show three consecutive traces of the centroid signal along the bunch while the instability is growing

HEADTAIL simulation, ξ_y=-0.2



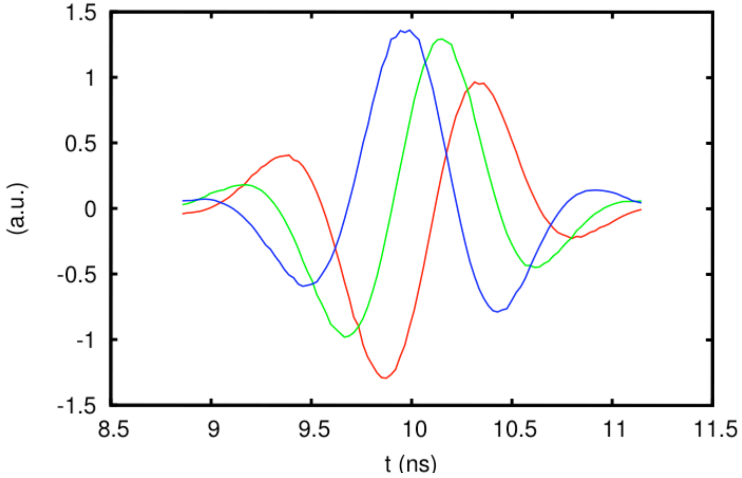
Measurement at the SPS (06.08.2007), ξ_y=-0.2



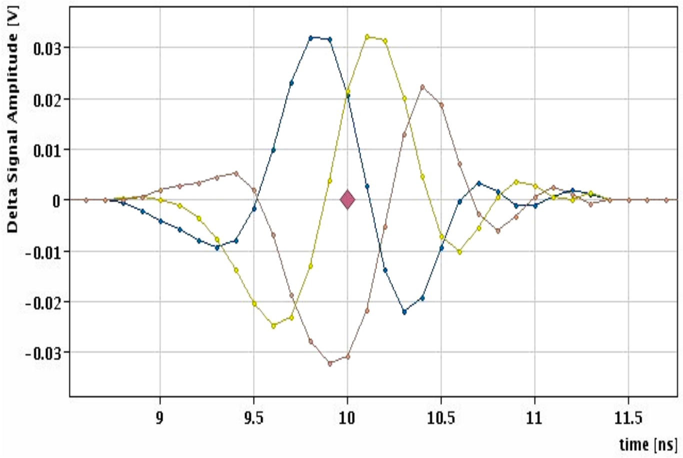
Example of simulation/measurements: **the head-tail instability**

- More benchmark of data and simulations for different values of chromaticity…

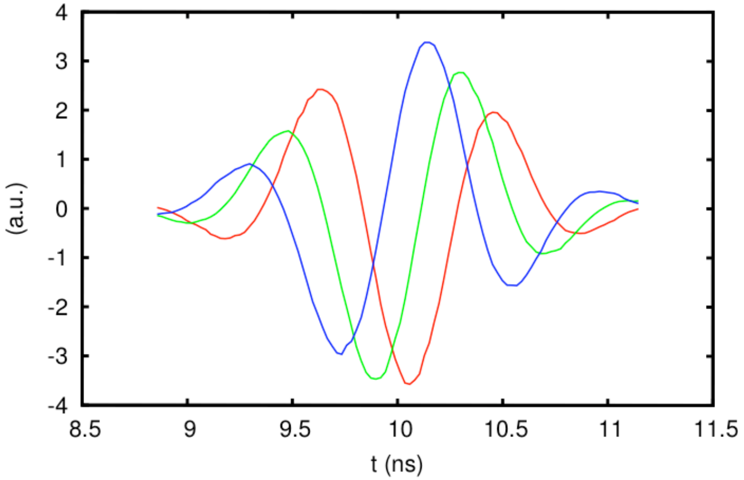
HEADTAIL simulation, ξ_y=-0.5



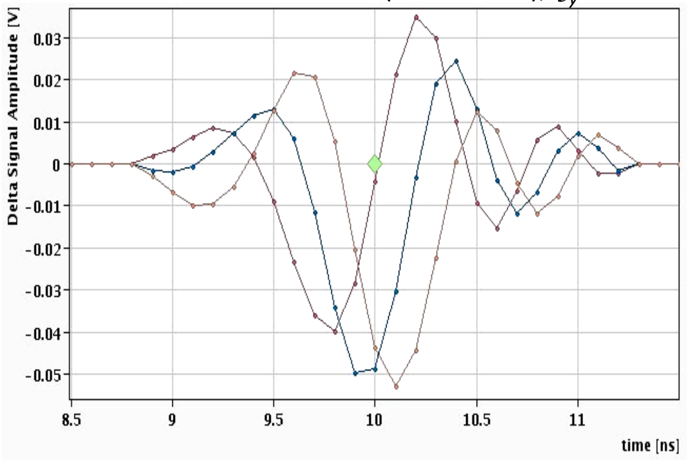
Measurement at the SPS (06.08.2007), ξ_y=*-0.5*



HEADTAIL simulation, ξ_y=-0.6



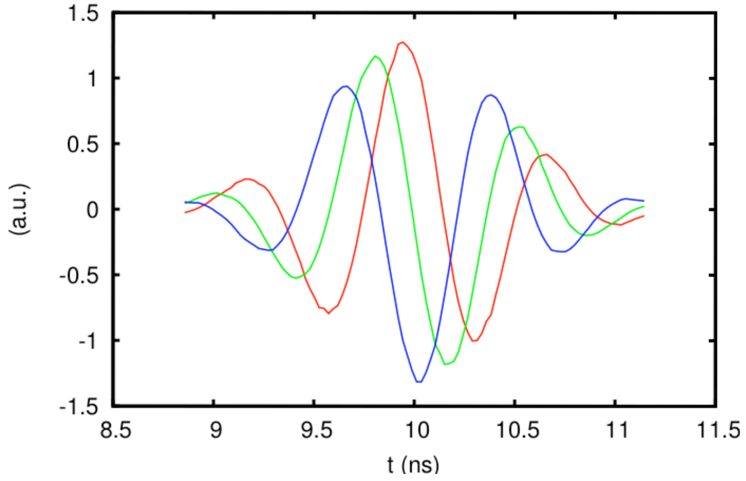
Measurement at the SPS (06.08.2007), ξ_y=*-0.6*



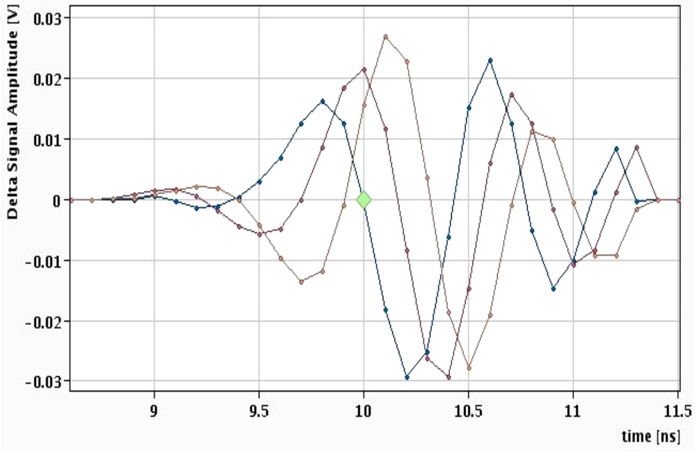
Example of simulation/measurements: **the head-tail instability**

- More benchmark of data and simulations for different values of chromaticity…

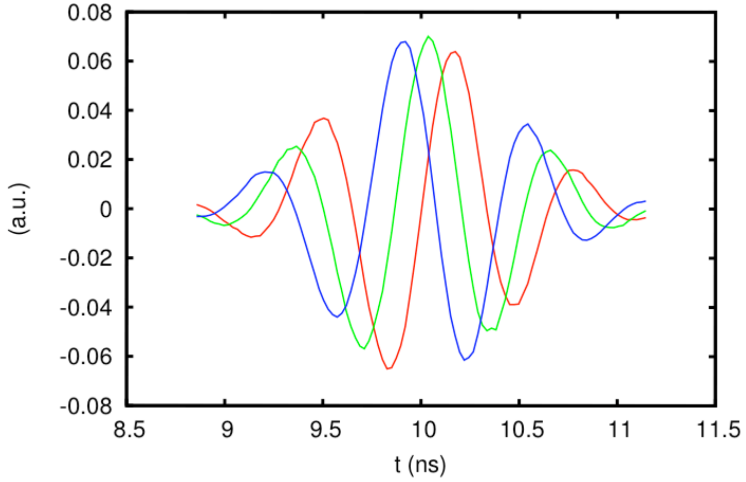
HEADTAIL simulation, ξ_y=-0.7



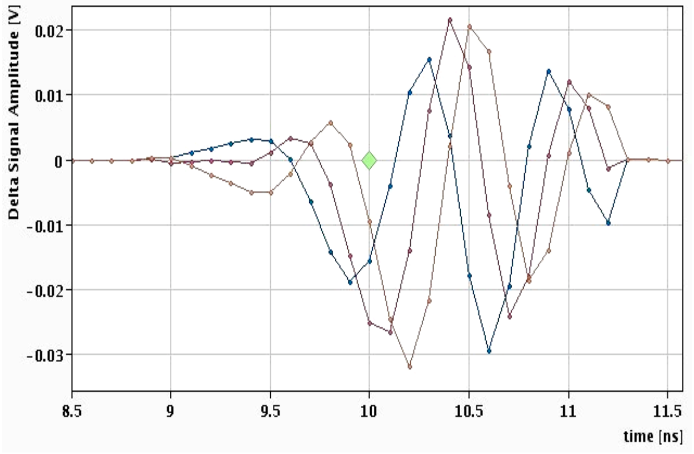
Measurement at the SPS (06.08.2007), ξ_y=-0.7



HEADTAIL simulation, ξ_y=-0.8



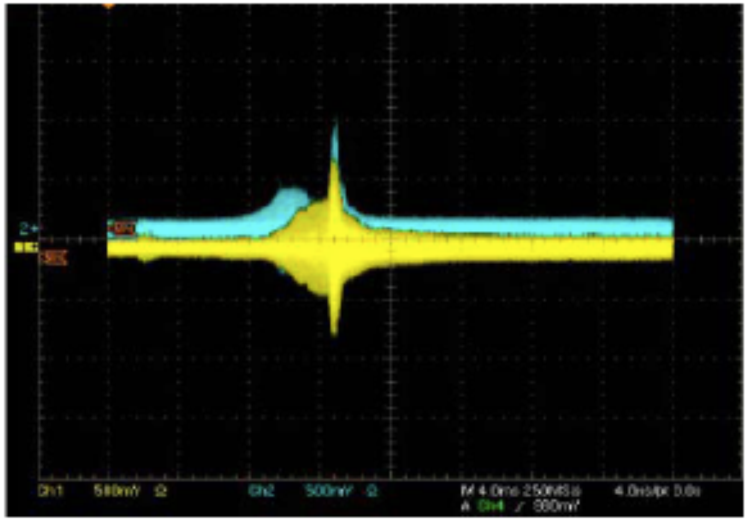
Measurement at the SPS (06.08.2007), ξ_y=-0.8



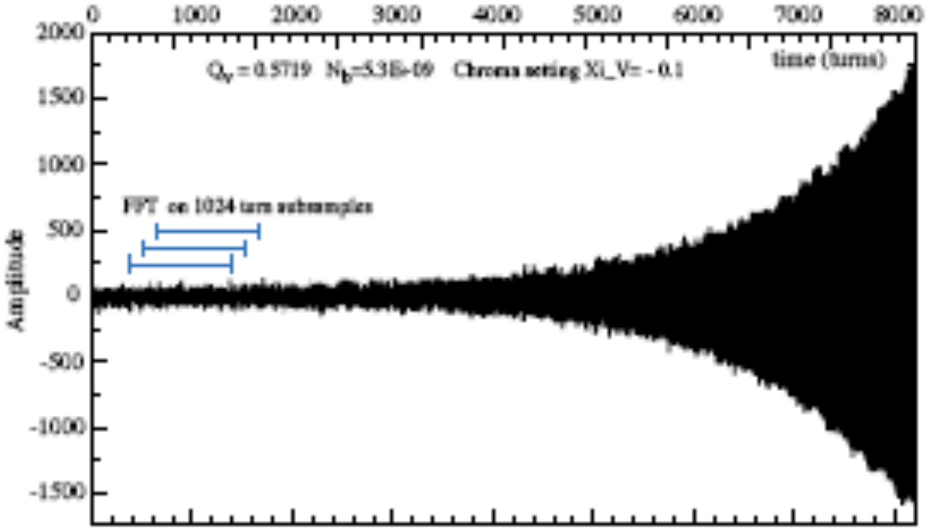
Example of measurements: **the head-tail instability**

- The growth rates of the head-tail modes are proportional to the real part of the machine impedance

  →The beam can be intentionally rendered unstable to obtain an estimation of the real part of the impedance of a machine by measuring the instability growth rate
  → If the bunch is long enough, the impedance spectrum can be probed by taking measurements at different chromaticity values.
  → Method applied to ORNL-SNS and to CERN-SPS



Single bunch instability measured at SNS
V. Danilov, et al., HB2006



Single bunch instability measured at SPS
H. Burkhardt et al. CERN-SL-2002-030

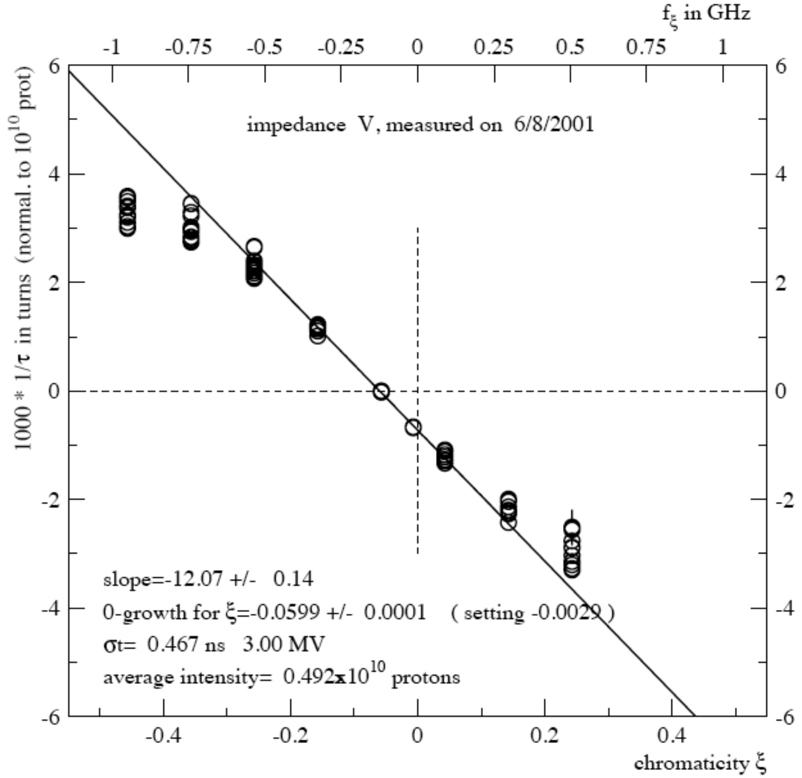
Example of measurements: **the head-tail instability**

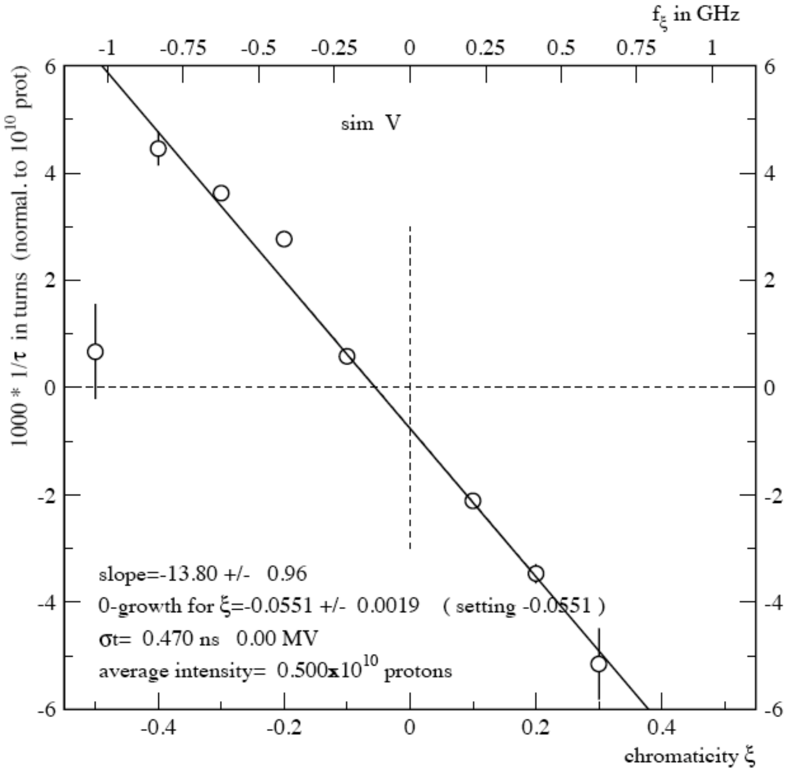
- The growth rates of the head-tail modes are proportional to the real part of the machine impedance

→ Growth/damping rates of the *l=0* mode are measured as a function of chromaticity

→ The bunch behavior is reproduced in simulation with a broad-band impedance model whose parameters are adjusted such as to match the observed trend

→ Example: SPS (2001)





Example of measurements: **the head-tail instability**

- Higher order head-tail modes (*l≥1*) are usually stabilized by tune spread and/or active feedback. However, if a high intensity beam stays in a machine long enough without sufficient tune spread and without feedback, these modes can also slowly grow.
- For example, a high intensity bunch becomes unstable in the CERN-PS over 1.2 s due to resistive wall

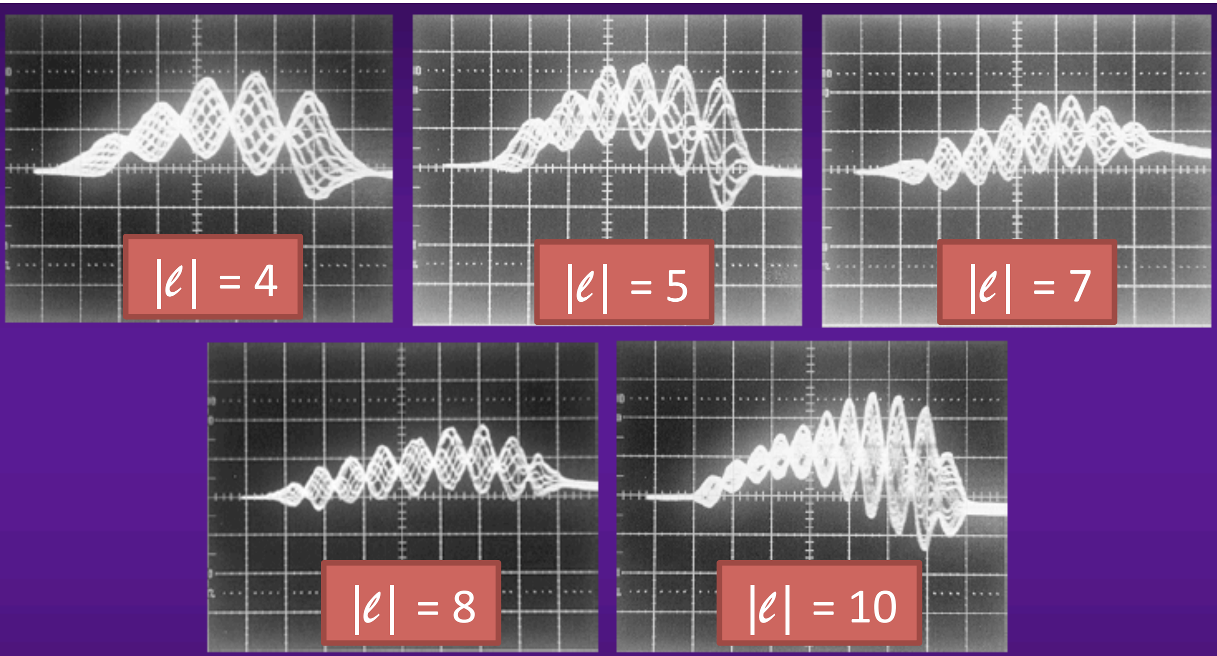


Figure 4: Measured ΔR signals from a radial beam-position monitor during 20 consecutive turns, in the PS with minimum coupling [5]: (a) $\xi_x \approx -0.5$, (b) $\xi_x \approx -0.7$, (c) $\xi_x \approx -1.1$, (d) $\xi_x \approx -1.2$, (e) $\xi_x \approx -1.3$. Time scale: 20 ns/div.

Example of simulation: **the head-tail instability**

- Higher order head-tail modes in the PS have also been simulated using the PS resistive wall impedance. These simulations are very demanding in terms of cpu time, because the bunch has to be tracked over about 500000 turns in order to see the effect arising from initial noise

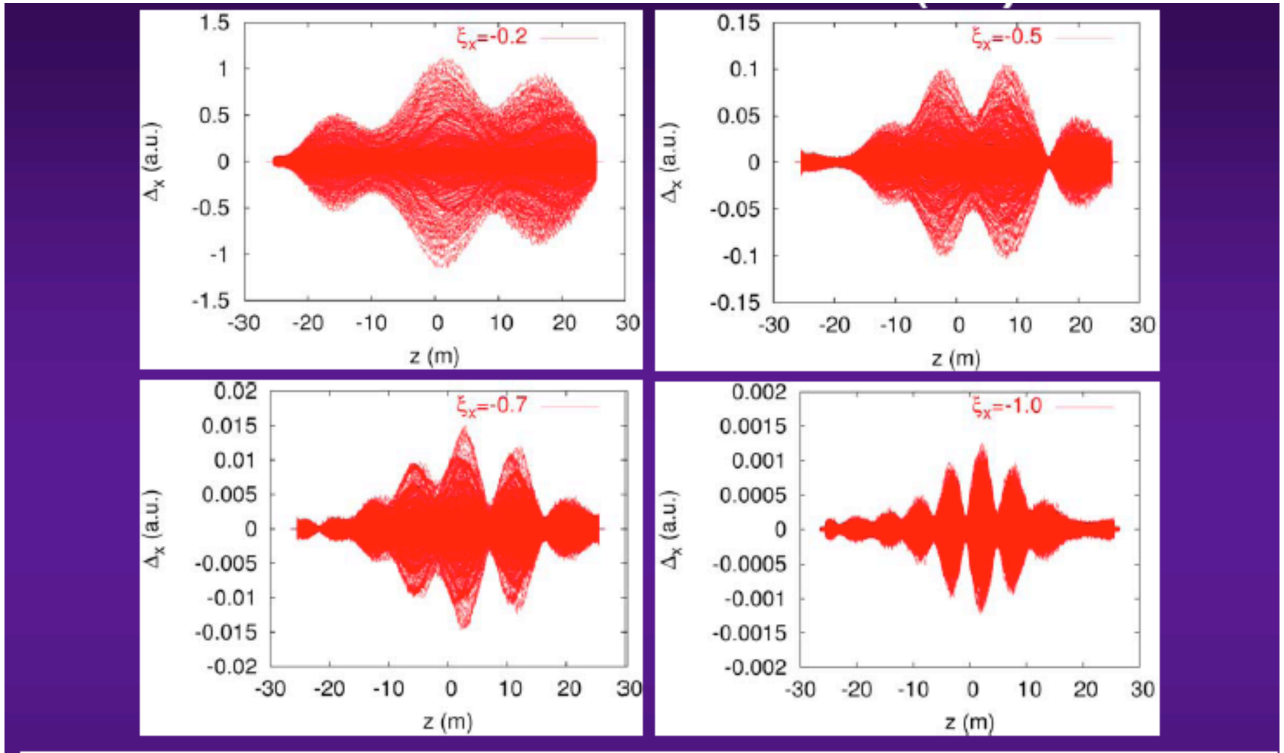


Figure 7: Examples of head-tail modes observed in the HEADTAIL simulations

Example of measurements/simulation: **the TMCI**

- The **Transverse Mode Coupling Instability** is another type of single bunch instability and has different features from the head-tail instability.

  ⇒ It does not depend on the chromaticity setting, and it actually occurs also for corrected chromaticity (in theory, for zero chromaticity)

  ⇒ It has a **threshold intensity** above which it appears.

  ⇒ The threshold value depends on the longitudinal emittance of the bunch, and bunches having lower longitudinal emittances tend to become more unstable

  ⇒ It is usually very fast (rise time shorter than the synchrotron period), that's why it is also called 'strong head-tail instability' or 'beam break-up'.

  ⇒ The shape of the Δ signal along the bunch is not caused by a head-tail phase shift from chromaticity, but depends on the spectrum of the driving impedance.

  ⇒ Mathematically, it appears when two head-tail modes merge at high intensity and two real solutions of the dispersion relation are replaced by a pair of complex conjugate solutions.

  ⇒ For many years the TMCI has been observed exclusively in lepton machines. The reason is that in hadron machines its threshold is increased by space charge and is usually higher than the threshold for the longitudinal microwave instability.
  However, the TMCI has been recently observed in the CERN-SPS (after the longitudinal impedance reduction campaign), in the CERN-PS and BNL-RHIC close to transition crossing.

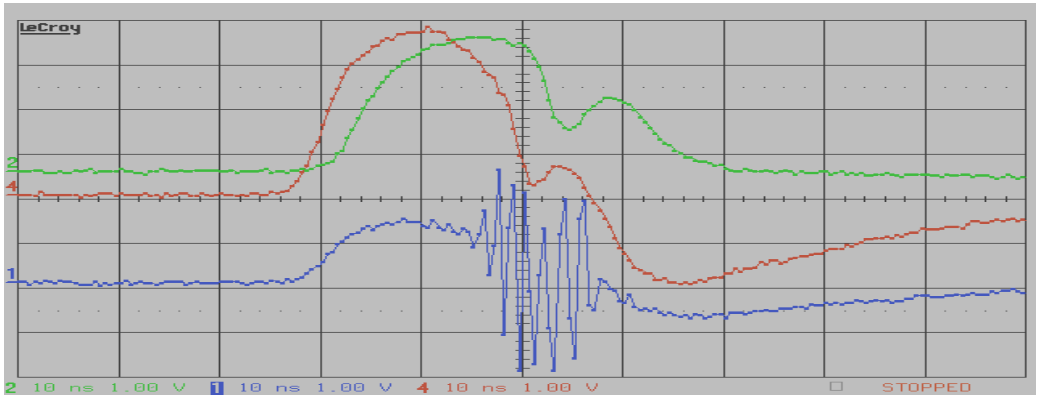
Example of measurements/simulation: **the TMCI**

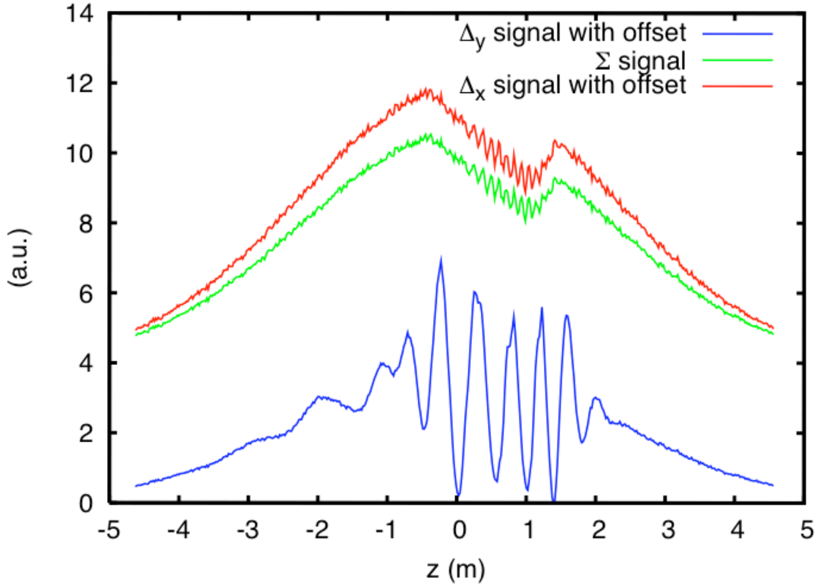
- The case of the PS high intensity bunch close to crossing transition energy

  ⇒ Beam loss was observed when crossing transition

  ⇒ The Δ_y signal along the bunch clearly showed turbulent vertical motion at a specific bunch location (i.e. a little off the peak towards the tail), where also the losses occurred

  ⇒ Simulations with a broad-band model could well reproduce the instability and the loss





Sum and Delta signals of the PS bunch at transition crossing.
Measurement (left) and simulation with a broad-band model (right)

Z_{eff}=3 MΩ/m @ 1 GHz

Movie

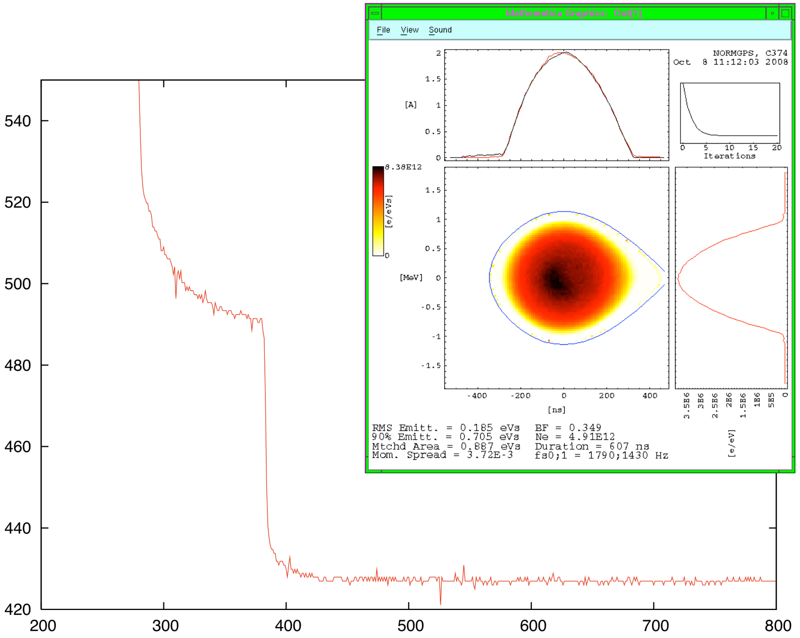
Example of experiment: the TMCI

- A PSB high intensity bunch becomes unstable along the ramp

  ⇒ Beam loss is observed at a specific point of the ramp when the damper is off

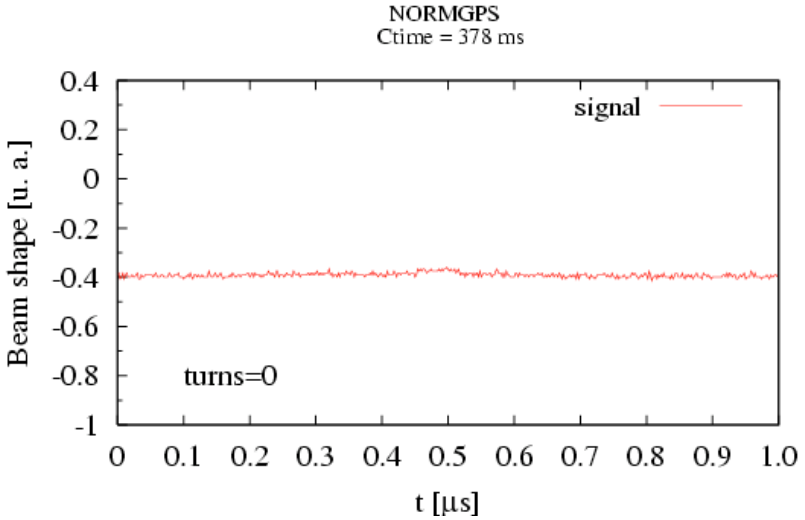
  ⇒ The Δ_x signal along the bunch clearly shows turbulent horizontal motion propagating from the tail of the bunch toward the head

  ⇒ Suspected TMCI



```
RMS Emitt. = 0.185 eVs   BF = 0.349
90% Emitt. = 0.705 eVs   Ne = 4.91E12
Mtchd Area = 0.887 eVs   Duration = 607 ns
Mom. Spread = 3.72E-3    fs0;1 = 1790;1430 Hz
```

Beam loss as measured by a Beam Current Transformer



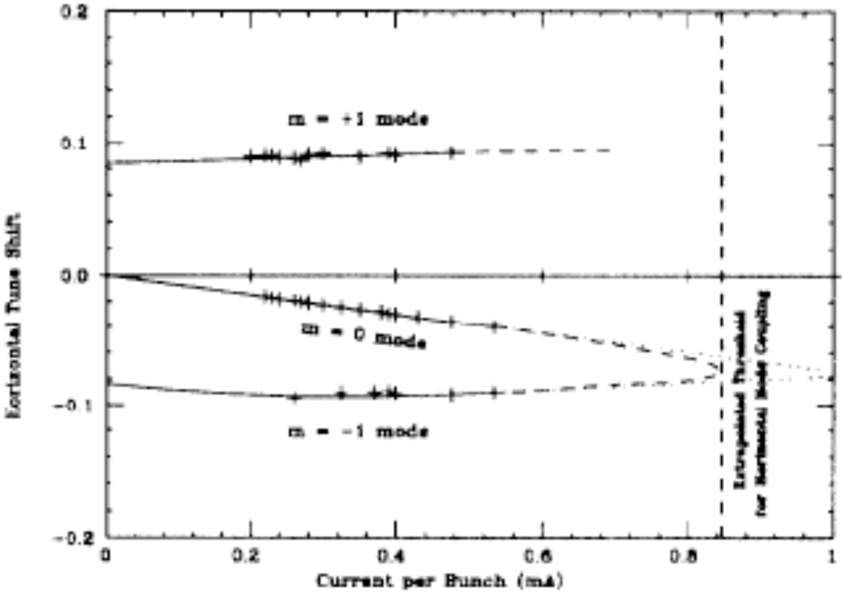
τ= 30 μs

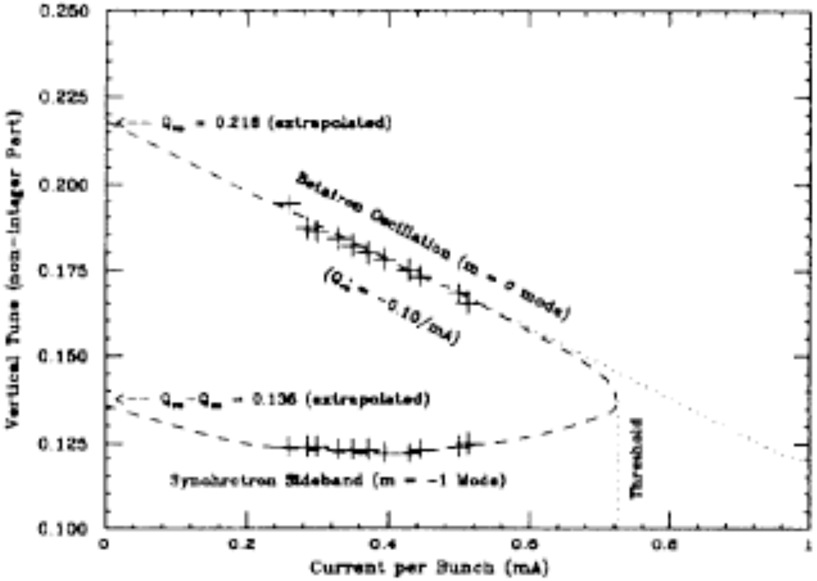
Example of measurements/calculation: Tune shift and TMCI

- Measurements of coherent tune shift as function of intensity in the CERN-LEP revealed other spectrum lines and in particular, the first synchrotron side bands (head-tail mode l=±1)

  ⇒ The two lines l=0 and l=-1 tend to merge as intensity increases

  ⇒ Measured values are in impressive agreement with the theoretical lines





B. Zotter, Comparison of Theory and Experiment on Beam Impedances: The Case of LEP, EPAC92

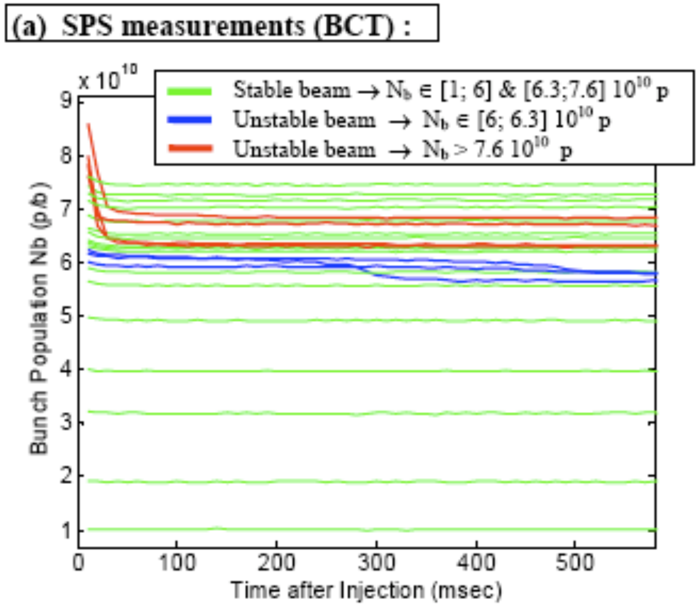
Example of measurements/simulation: Tune shift and TMCI

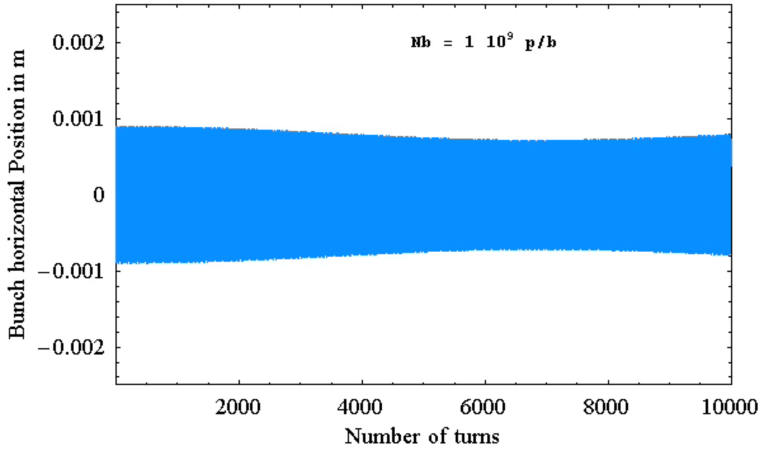
- Measurements of coherent tune shift as function of intensity in the SPS have revealed that, using a low longitudinal emittance bunch, a vertical TMCI can be observed at injection above a certain intensity threshold

⇒ Beam loss is observed at injection in some intensity ranges

⇒ The Δ_y signal along the bunch clearly shows turbulent vertical motion propagating from the tail of the bunch toward the head

⇒ A moderately unstable intensity range seems to be followed by a stable one before getting into a strong instability region



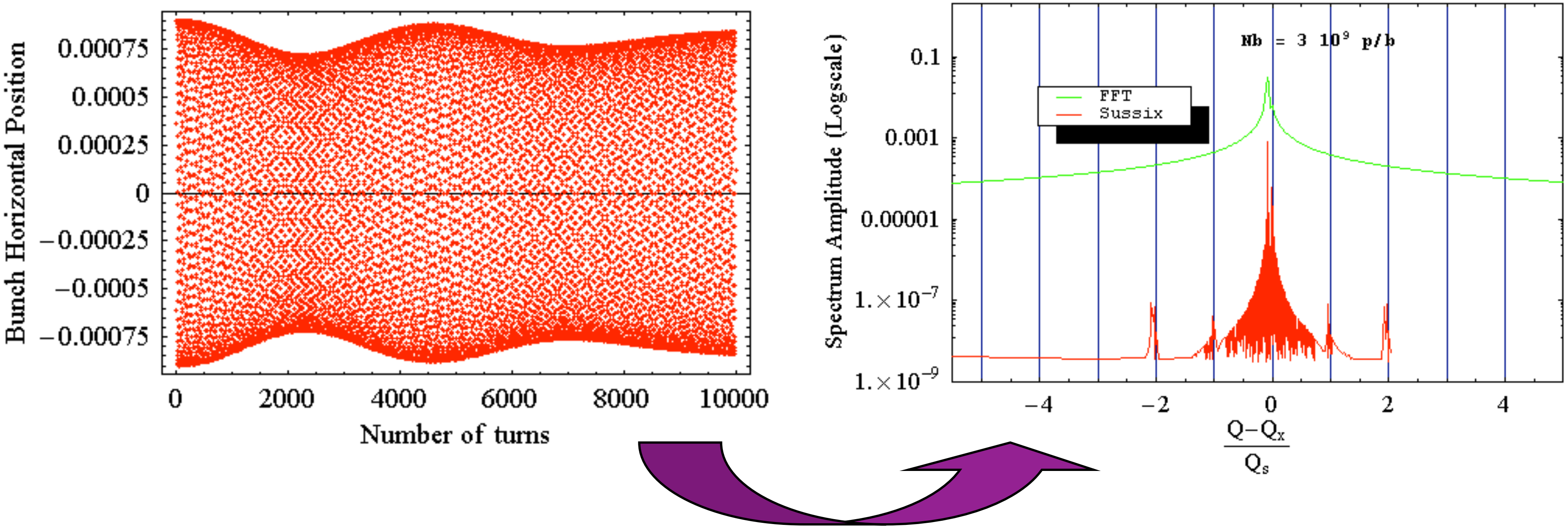


The simulated evolution of the bunch predicted the existence of slightly unstable regions for intensities lower than 8 x 10^{10}

Movie

Benchmarking MOSES and HEADTAIL (I)

The fine structure of the coherent modes can be revealed from the HEADTAIL output of the centroid motion by applying SUSSIX to the complex BPM signal $(x + j\,\beta_x x')$

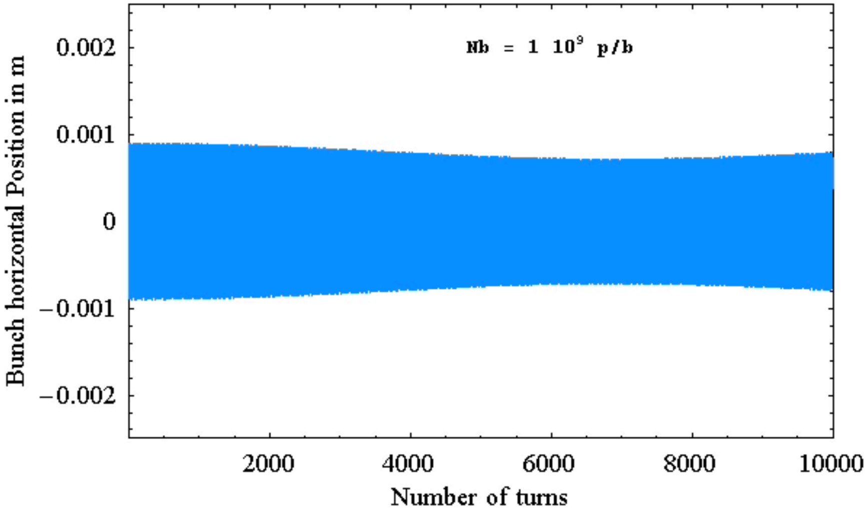


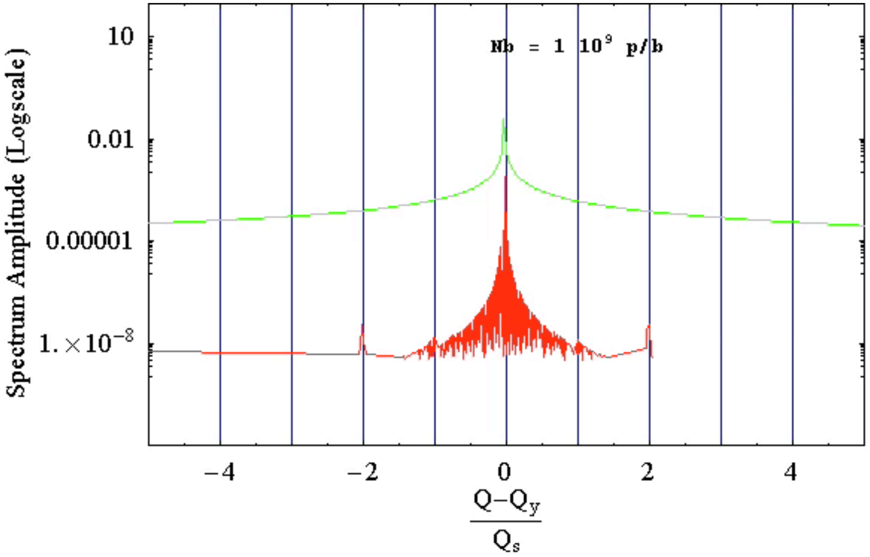
Standard FFT or SUSSIX

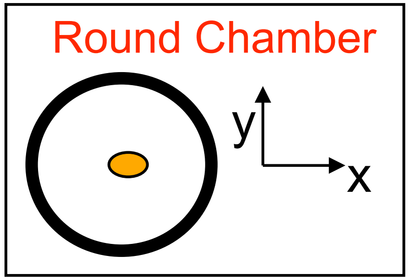
(*) *Sussix code* : R. Bartolini, F. Schmidt*, SL Note 98-017AP, CERN 1998*
Theory behind Sussix : R. Bartolini, F. Schmidt, *LHC Project Report 132, CERN 1997*
J. Laskar et al.*, Physica D 56, pp. 253-269 (1992)*

Benchmarking MOSES and HEADTAIL (II)

Scanning in intensity we can observe how the mode in the spectrum shift....

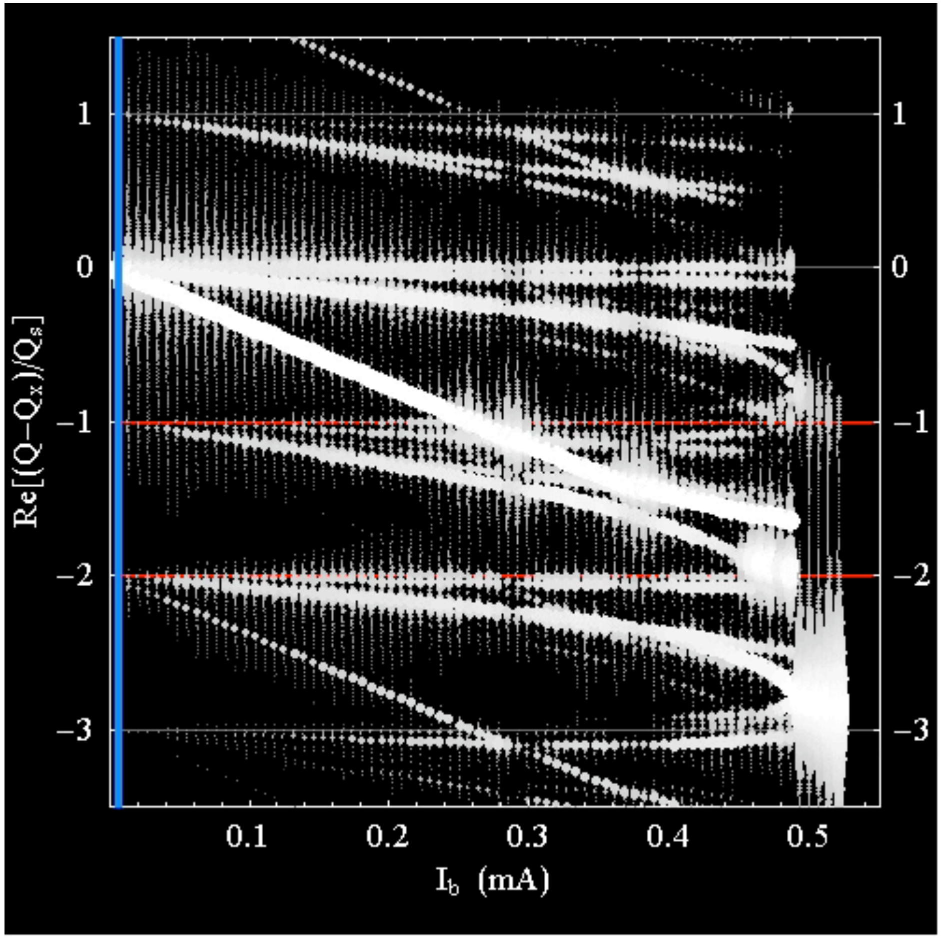


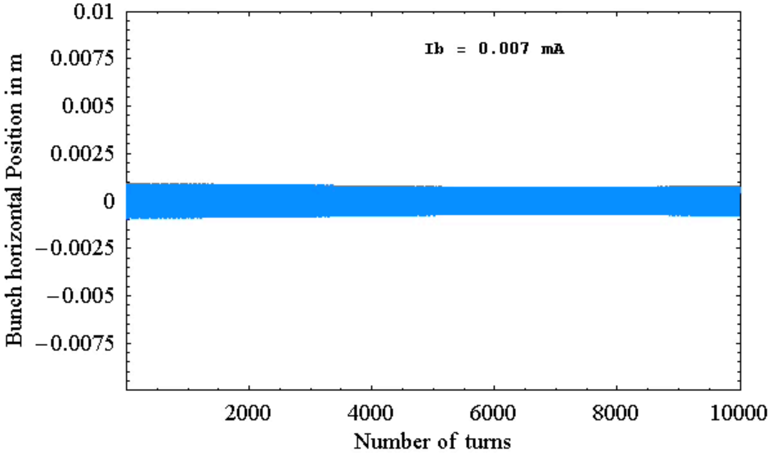




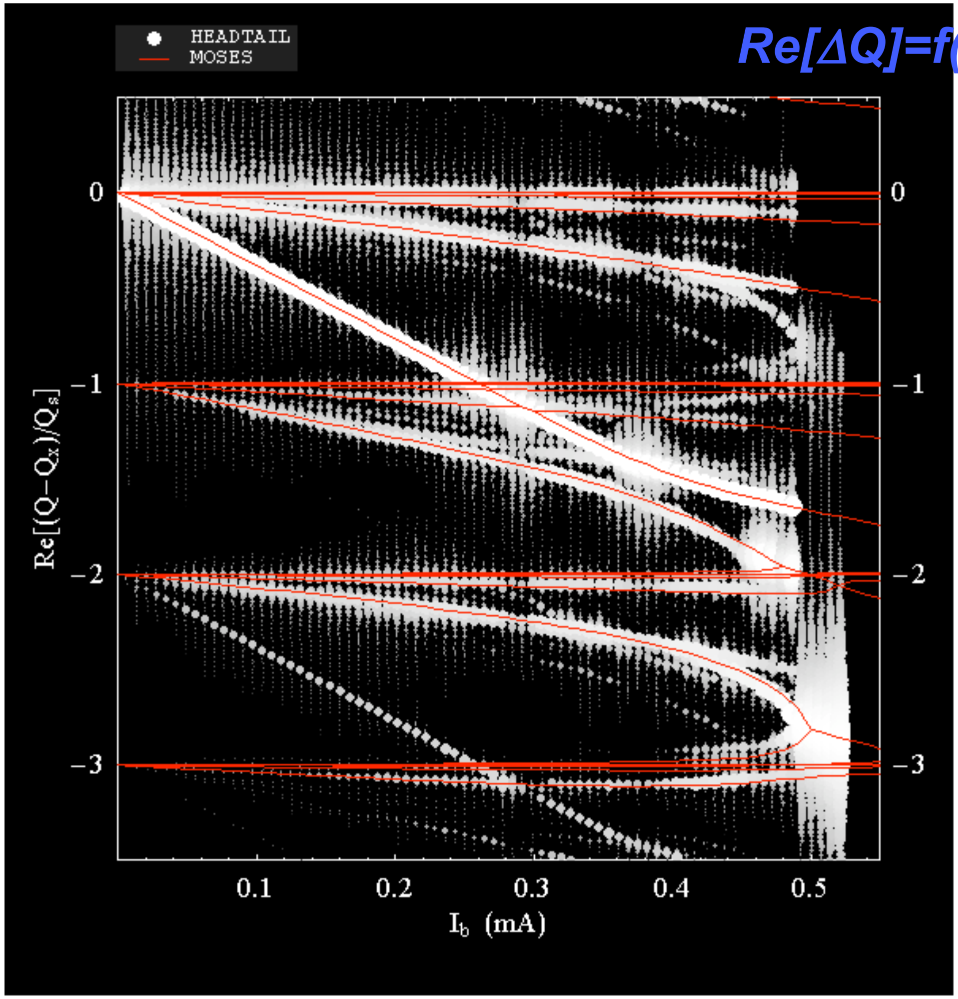
Benchmarking MOSES and HEADTAIL (III)

Round beam pipe / no chromaticity / no coupling - displaying ***Re[ΔQ]=f(I_b)***





Benchmarking MOSES and HEADTAIL (iV)



$\mathrm{Re}[\Delta Q]=f(I_b)$ and comparison with MOSES

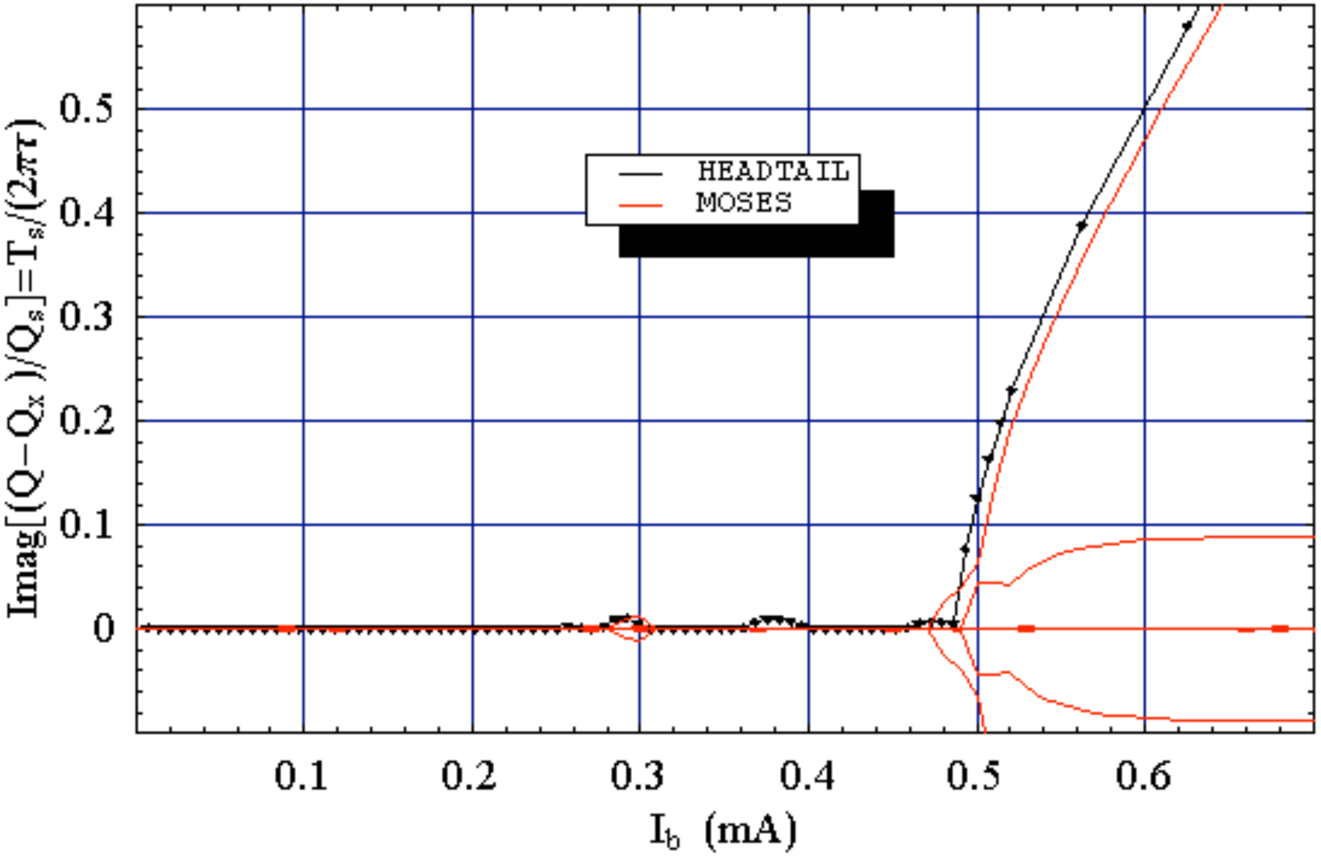
Results from **MOSES** coherent mode analysis and **HEADTAIL** are superimposed and the agreement is excellent

⇒ Most of the radial modes per azimuthal number can be seen from HEADTAIL

⇒ There are few *ghost* lines, probably due to a small initial mismatch of the bunch in the bucket

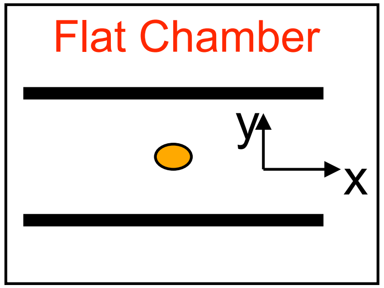
Benchmarking MOSES and HEADTAIL (V)

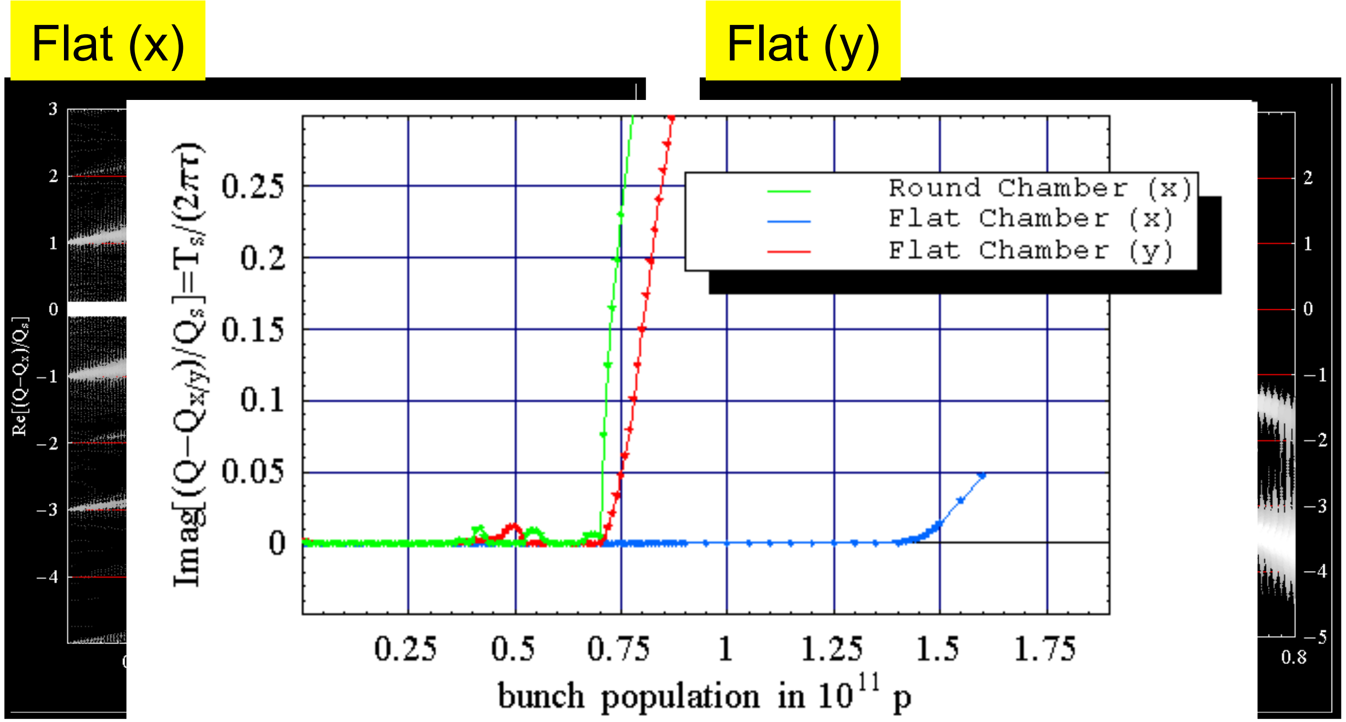
Im[ΔQ]=f(I_b) and comparison with MOSES



Also the agreement between the predicted instability growth rates is very good

HEADTAIL with flat pipe





Mode analysis with flat pipe:

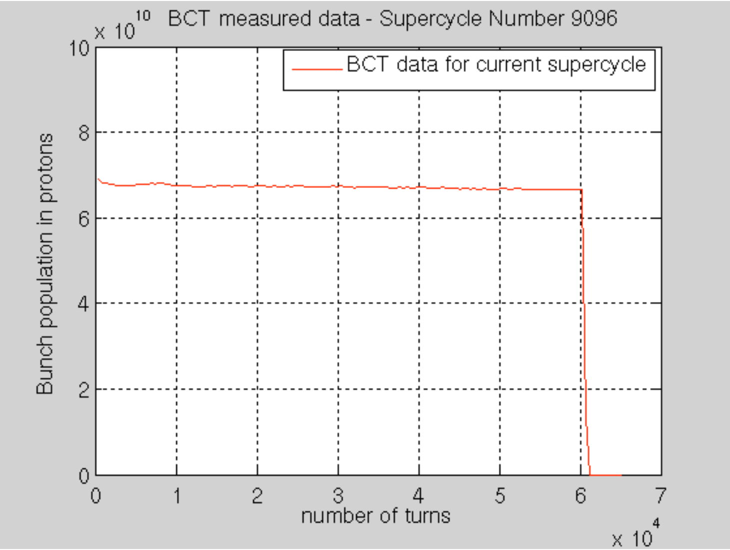
→ Modes shift differently in x and y

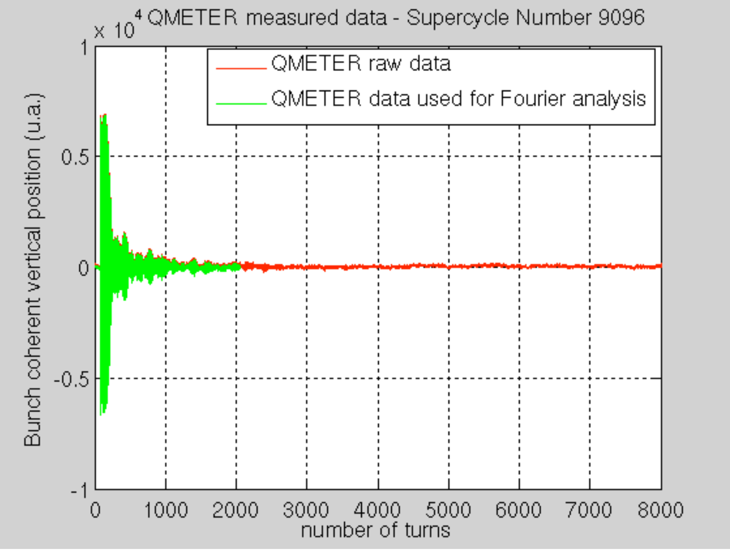
→ The instability threshold in ***x*** is found to be about twice the threshold in ***y***

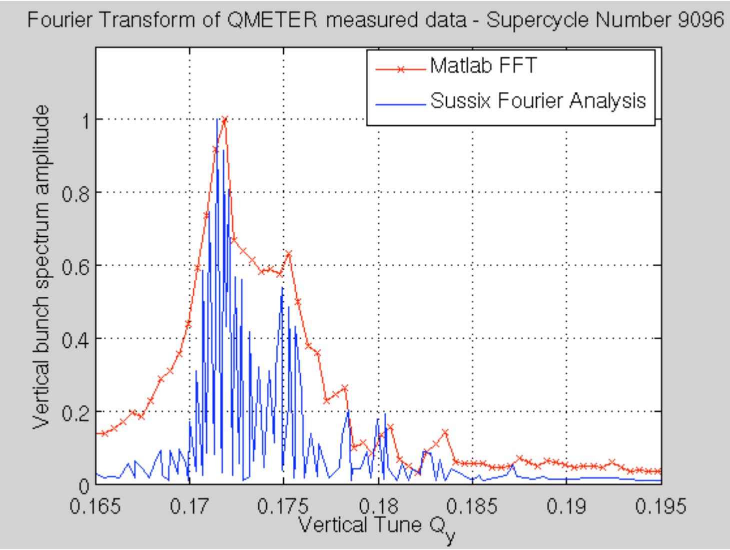
→ This suggests that linear coupling could help in this case....

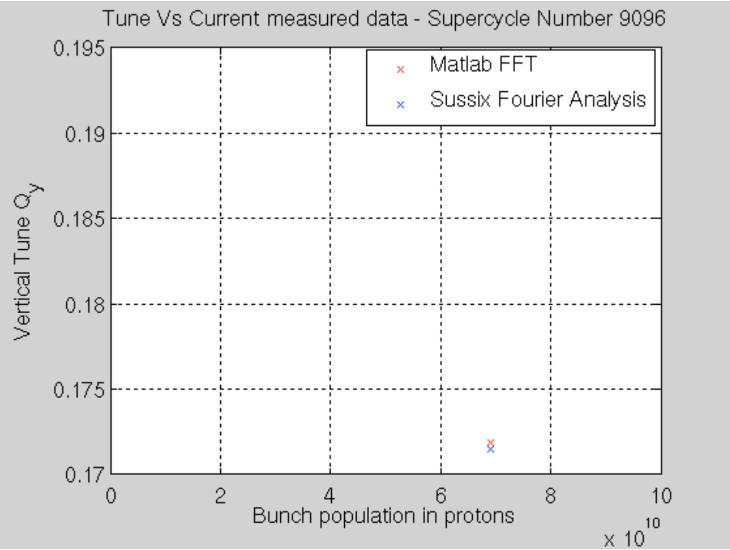
Example of measurements: Tune shift

- What we can measure below the TMCI threshold….









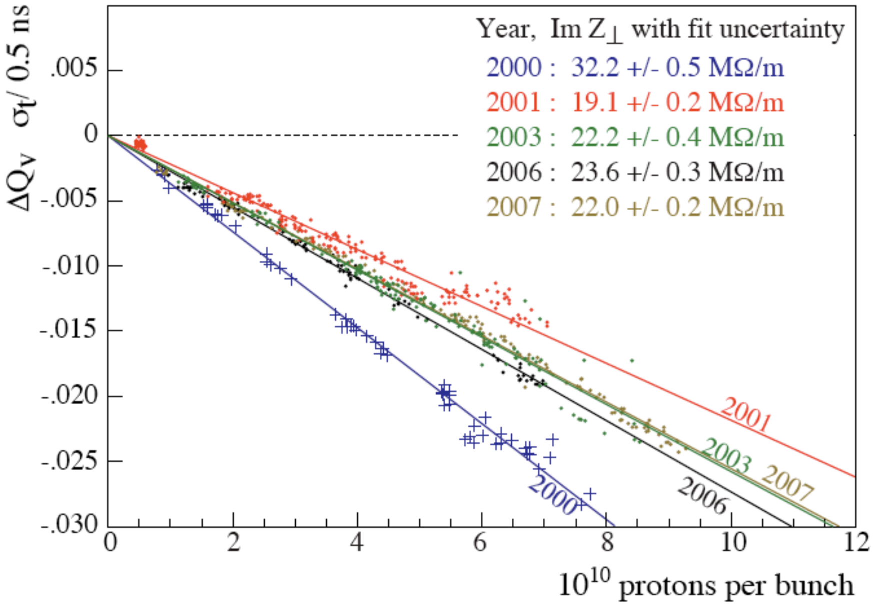
Example of measurements: **Tune shift**

- Measurements of coherent tune shift as function of intensity in the CERN-SPS

  ⇒ From the slope of the tune shift one can infer the low frequency imaginary part of the machine impedance (iZ_{eff}). Machines with flat beam pipes show usually no tune shift in the horizontal plane and significant tune shift in the vertical plane

  ⇒ Tune shift measurements done with high longitudinal emittance bunches can extend to high intensities because the TMCI threshold is higher

Vertical coherent tune shift with intensity at 26 GeV, scaled to 0.5 ns

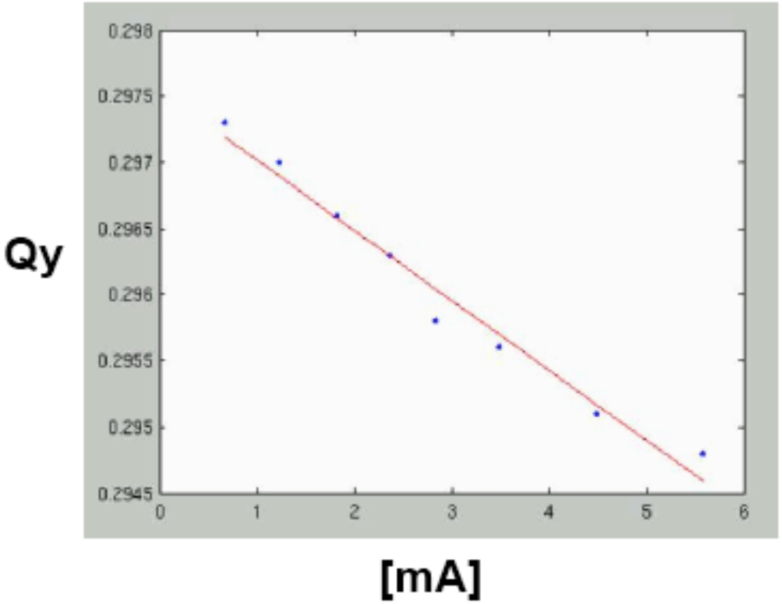


Example of measurements: Tune shift

- Measurements of coherent tune shift as function of intensity at the SSRF (Shangai Synchrotron Radiation Facility)

  ⇒ J. Bocheng, C. Guanglin, C. Jianhui, "Collective effects of SRRF storage ring 3 GeV Phase I commissioning", SSRF internal note, April 2008; J. Bocheng, "Impedance budget of SSRF storage ring", SSRF internal note, April 2008.

- Vertical: $(Z_\perp)_{eff}$ = **98 ~ 136 kΩ/m** measured from the coherent tune shift, which is nearly a **factor of 2** above expectation.



- $(I_{th})_{RW}$ ~64 mA (ξ_y = 0.1) and > 100 mA (ξ_y > 0.5).

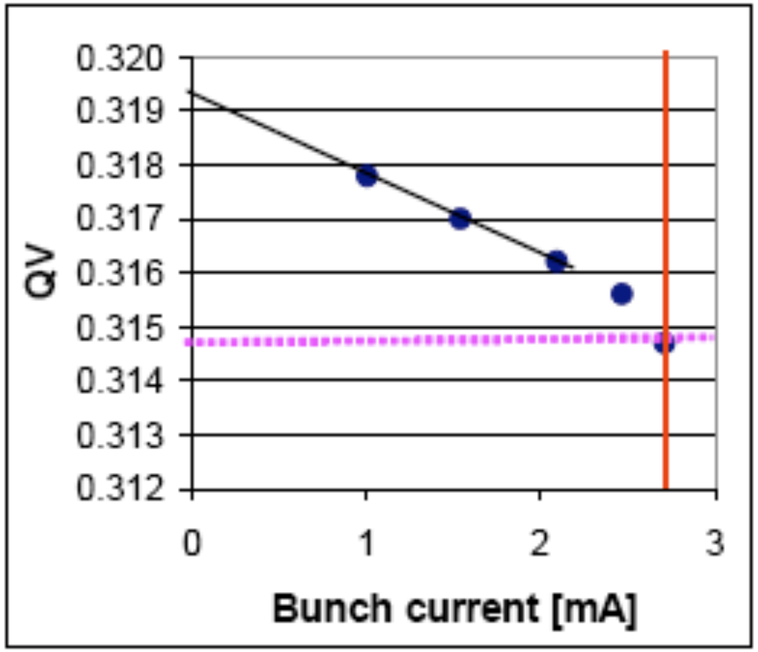
- Ion instabilities disappeared 1 month after the start of commissioning when the vacuum improved to 5×10^{-10} Torr.

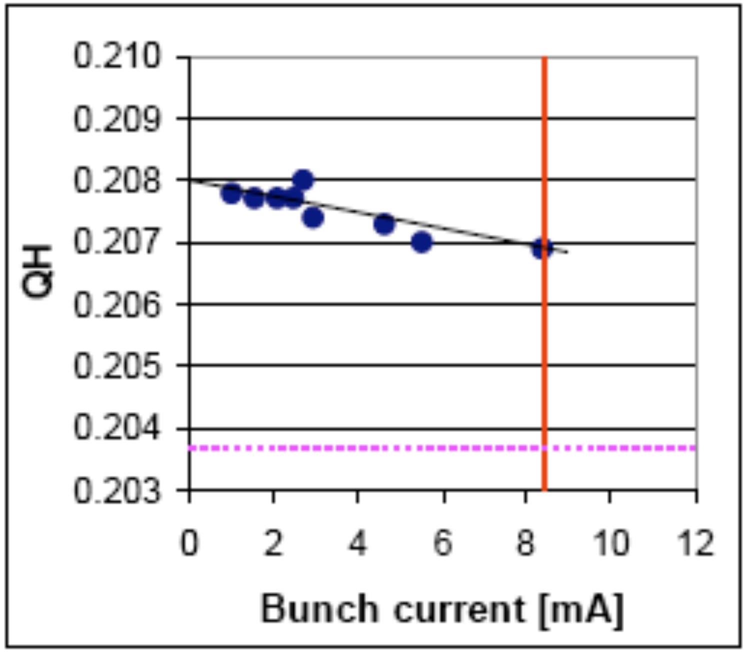
Example of measurements: **Tune shift**

- Measurements of coherent tune shift as function of intensity at the Soleil

  ⇒ R. Nagaoka, MP. Level, L. Cassinari, ME. Couprie, M. Labat, C. Mariette, A. Rodriguez, R. Sreedharan, PAC07

  ⇒ Measured Z_{eff} is measured to be larger than expected by a factor of ~2 both in H and V planes.



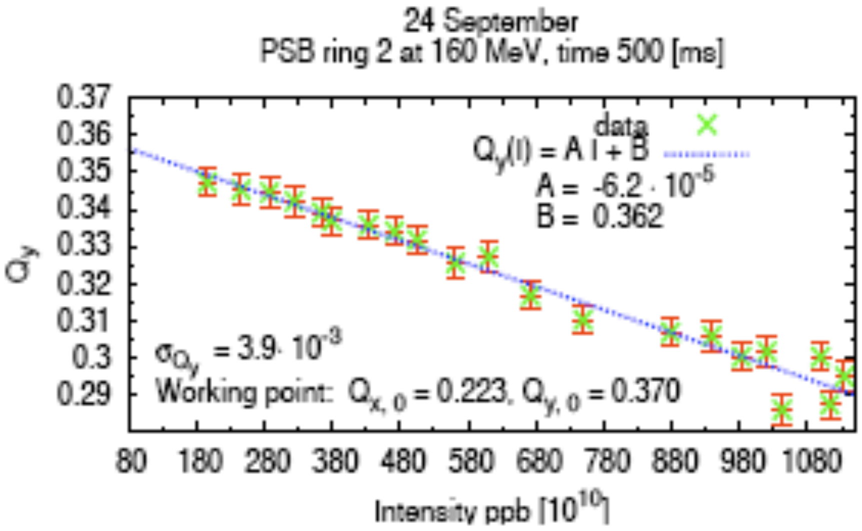


Example of measurements: Tune shift

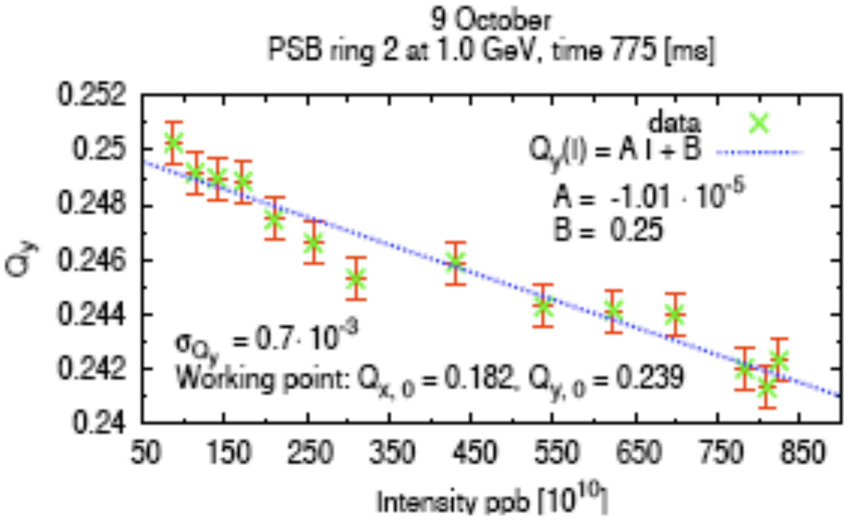
- Measurements of coherent tune shift as function of intensity in low energy machines is more tricky because the contribution of the beam images (indirect **S**pace **C**harge) has to be disentangled from the contribution of the **M**achine **I**mpedance (in principle independent of energy)

  ⇒ Measurements at different energies can be used for this purpose

  ⇒ The method has been applied recently to the CERN-PSB



Z_{eff}=14 MΩ/m=Z_{MI}+Z_{SC}(160 MeV)



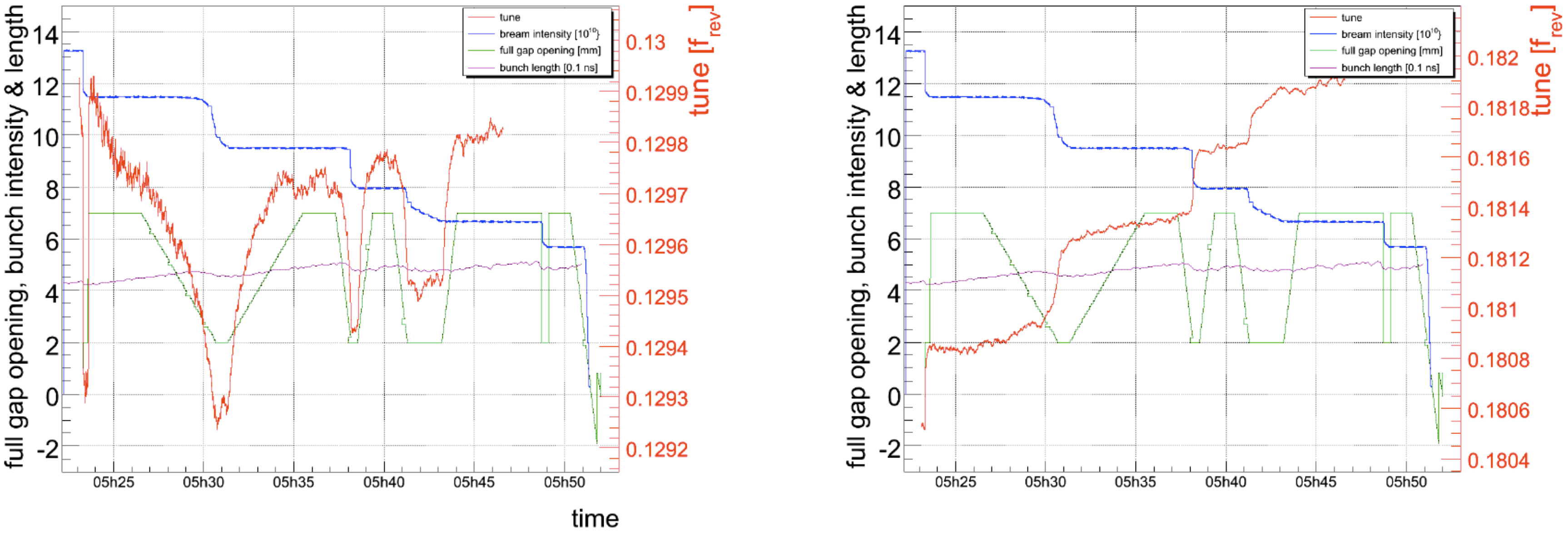
Z_{eff}=5 MΩ/m=Z_{MI}+Z_{SC}(1 GeV)

Example of measurements: Tune shift

- Some times the tune shift can be measured changing in a controlled way a known impedance source inside the machine

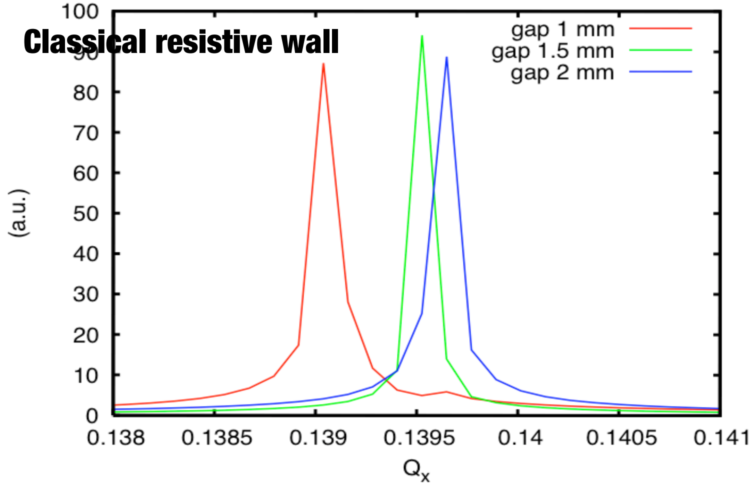
  ⇒ Typical "tunable" impedance sources are movable collimators, scrapers or other intercepting devices, as the transverse impedance scales like g^{-3} (g being the device gap)

  ⇒ Tune measurement in the CERN-SPS while a prototype of LHC collimator (installed in the machine for test purposes) was being moved inward and outward in the horizontal plane. The vertical tune variation is due to the beam loss caused by the collimator when moved in

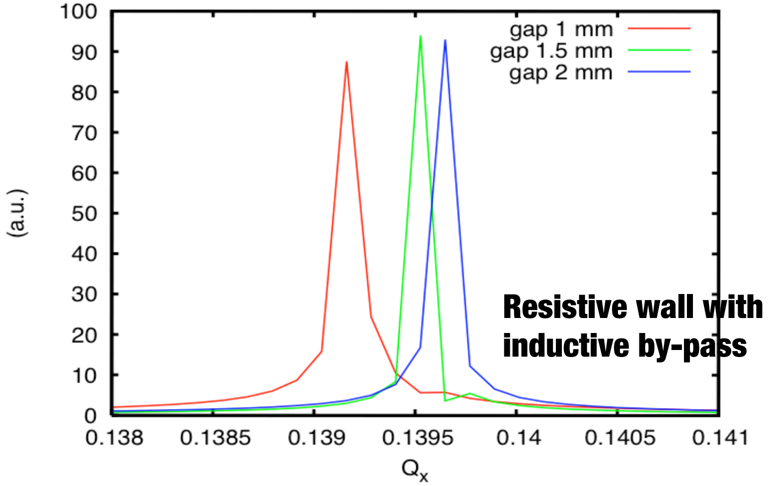


Collimator MD@SPS on the 1 November 2006

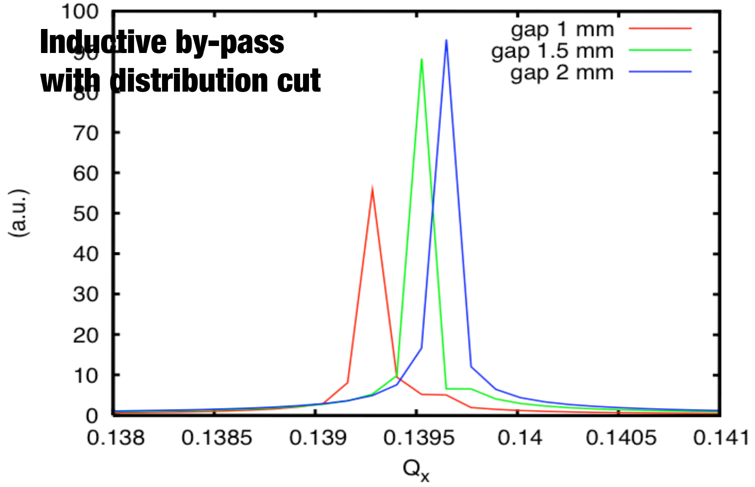
Simulations of the collimator with HEADTAIL



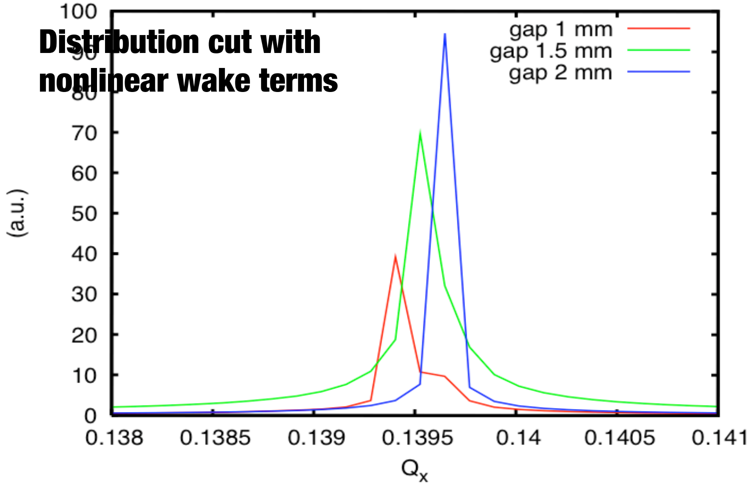
Classical resistive wall



Resistive wall with inductive by-pass

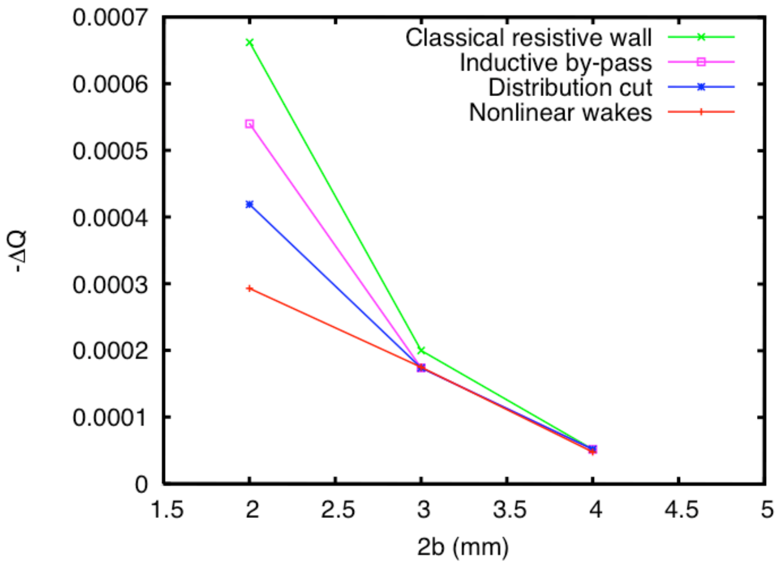


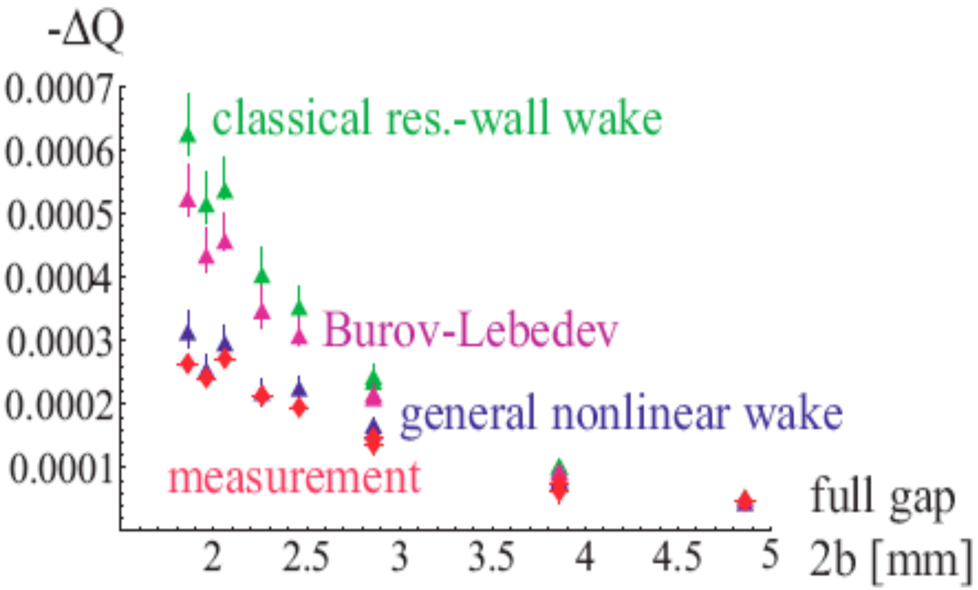
Inductive by-pass with distribution cut



Distribution cut with nonlinear wake terms

Simulations of the collimator with HEADTAIL





Comparing the tune shifts extrapolated from **HEADTAIL simulations** (left plot) with the experimental ones (right plot, **red** points) and those from analytical theory (right plot, **green**, **magenta**, **blue** points), the agreement is excellent.

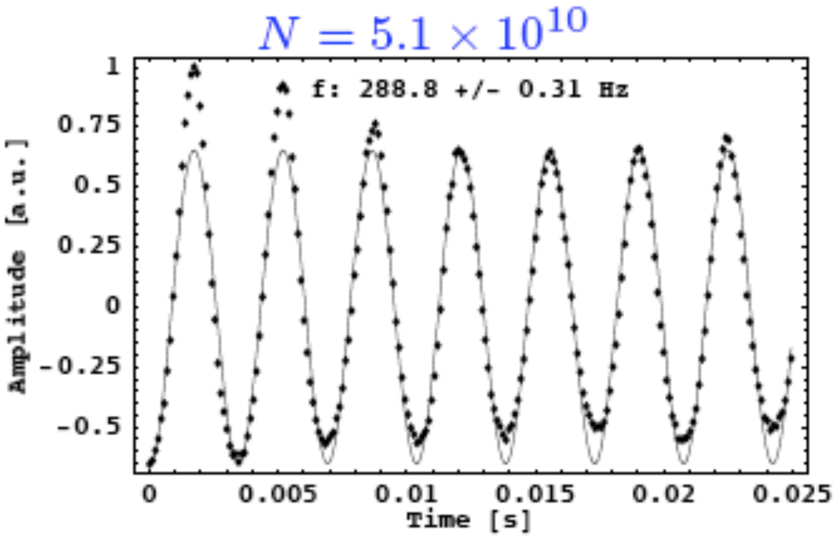
Example of measurements: Tune shift (longitudinal)

- Measurements of synchrotron tune shift as function of intensity can be also done in the longitudinal plane in order to estimate the longitudinal impedance

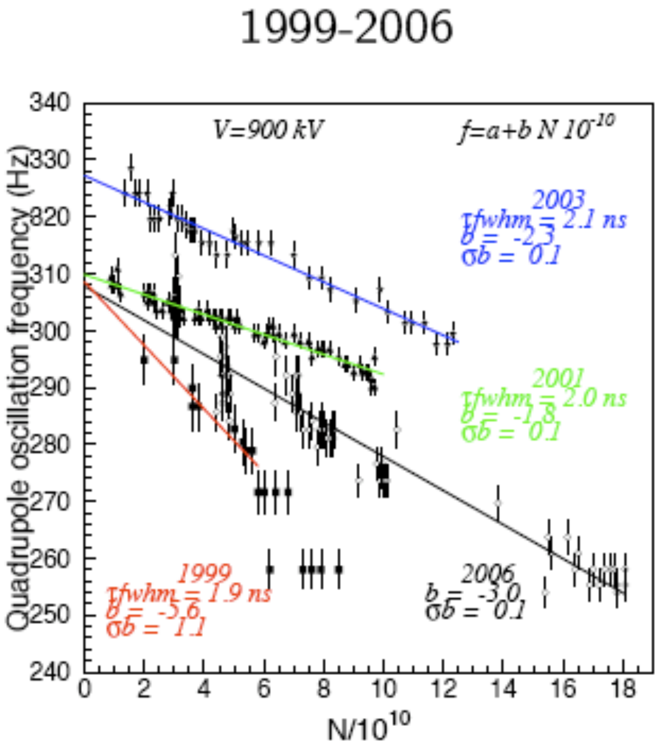
  ⇒ The shift appears in the quadrupole mode, therefore the technique uses e.g. the synchrotron oscillations of a bunch injected with a mismatch

  ⇒ Q_s can be extrapolated from bunch length or peak amplitude measurements

  ⇒ Example: SPS measurements by E. Shaposhnikova, T. Bohl, J. Tuckmantel



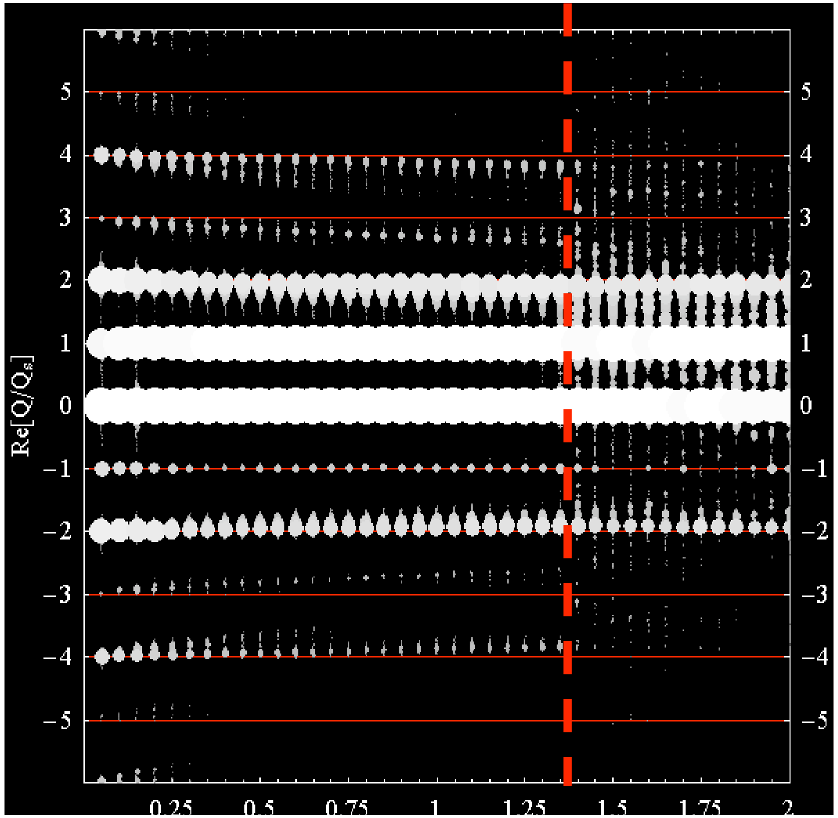
$Z_{\parallel eff}/n \approx 5\ \Omega$



Example of simulations: **Longitudinal impedance acting on an SPS bunch**

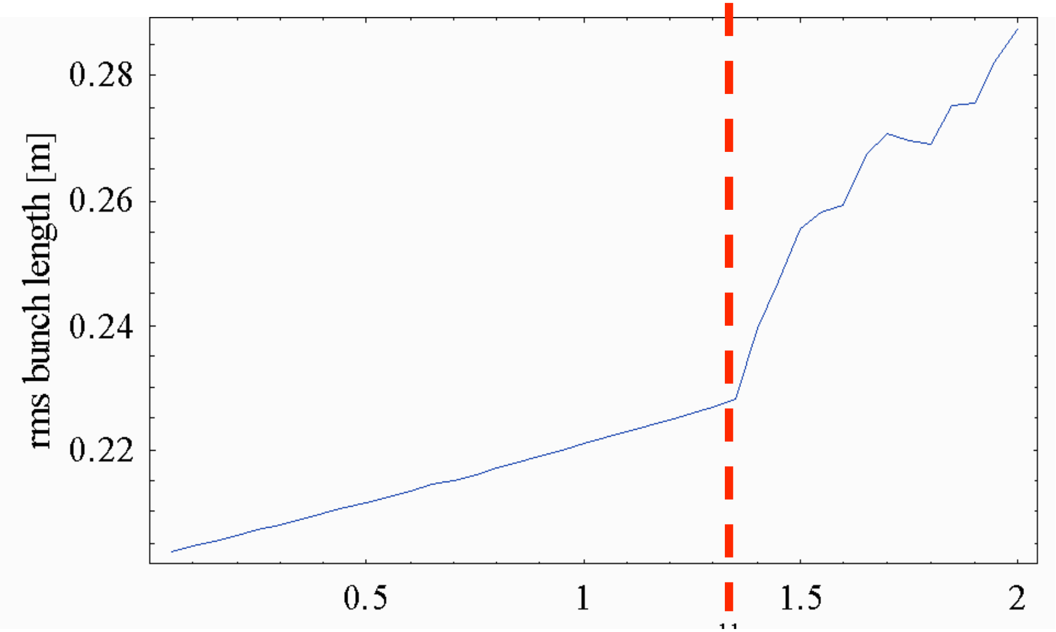
- Simulating the effect of a longitudinal impedance on an SPS bunch we can clearly distinguish the effects in lower and higher intensity regimes

  ⇒ Potential well distortion regime shows with a linear increase of the bunch length as a function of the bunch intensity, while the unstable regime is characterized by a change of slope in bunch lengthening



Potential Well Bunch Lengthening regime

Microwave Instability regime



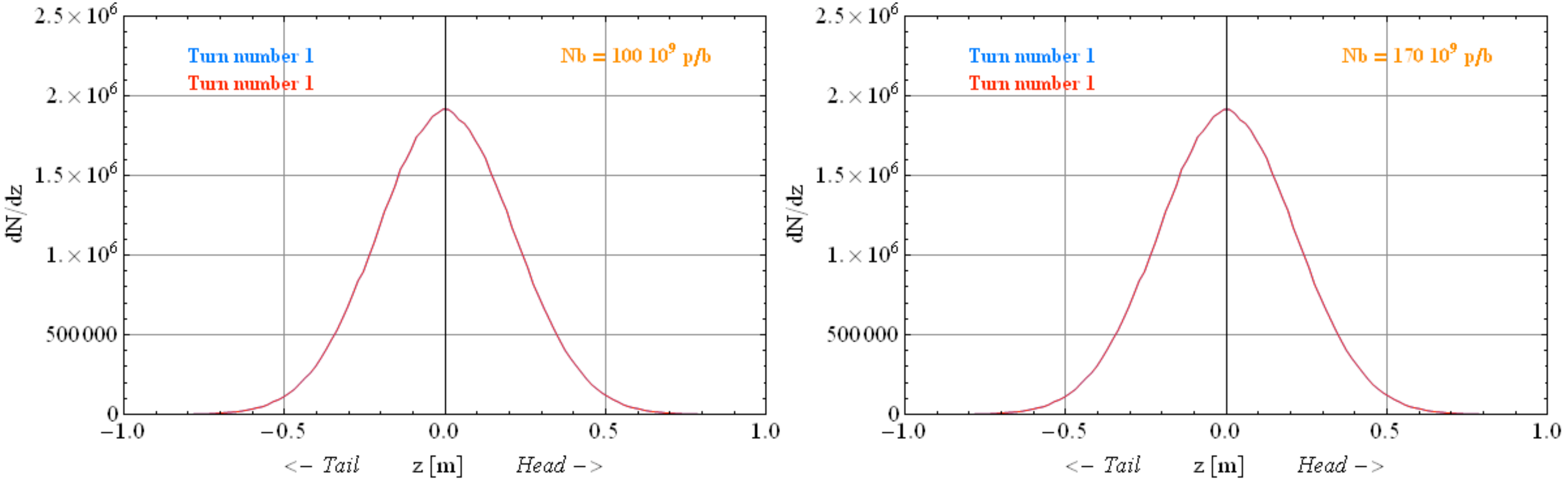
Broad-band, $Z_{//}/n$=10 Ω, f_r=700 MHz

Example of simulations: Longitudinal impedance acting on an SPS bunch

- Simulating the effect of a longitudinal impedance on an SPS bunch we can clearly distinguish the effects in lower and higher intensity

  ⇒ Bunch lengthening regime: slow evolution towards a new equilibrium with a slightly shifted synchronous phase due to energy loss.

  ⇒ Unstable regime: micro-bunching appear.



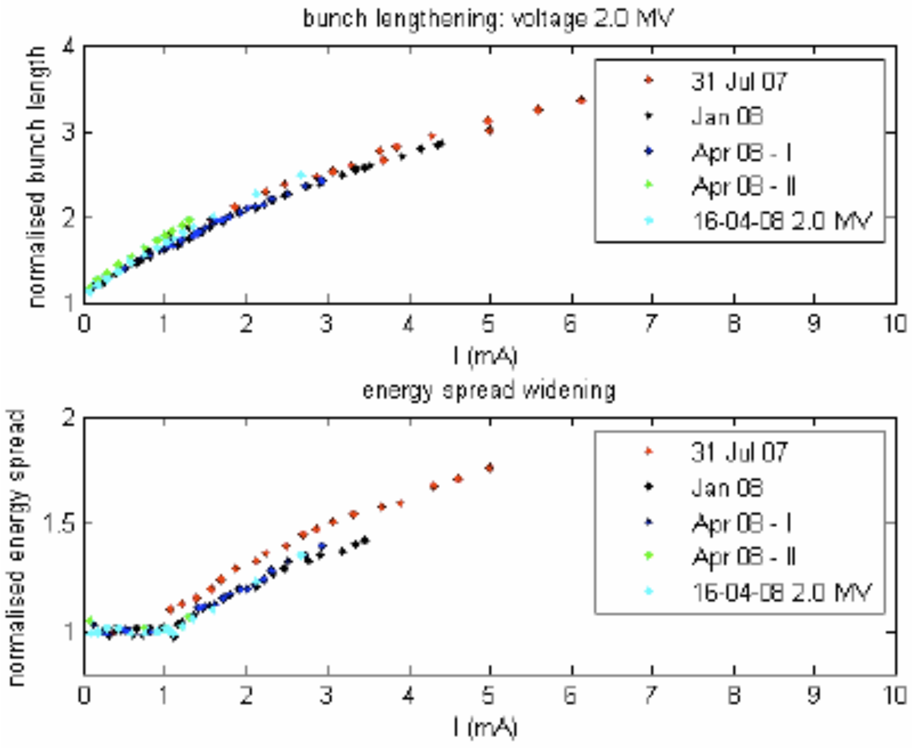
Bunch shape evolution in the regime of potential well distortion (10^{11} ppb, left movie) and just above the threshold for microwave instability (1.7×10^{11} ppb, right movie)

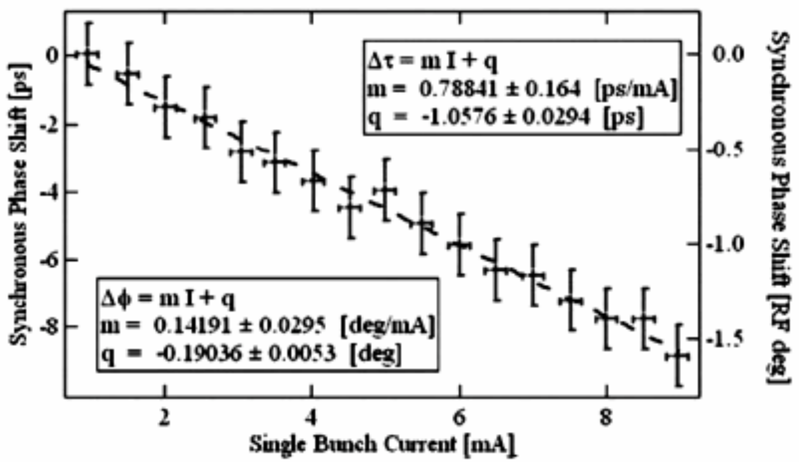
Example of measurements: **Other methods to estimate $Z_{\parallel eff}$**

- In order to estimate the longitudinal impedance, it is also possible to look at

  ⇒ Bunch lengthening (ex. DIAMOND, R. Bartolini)

  ⇒ The energy loss measured through the synchronous phase shift (ex. Australian light source, R. Dowd, M. Boland, G. LeBlanc, M. Spencer, Y. Tan, PAC07





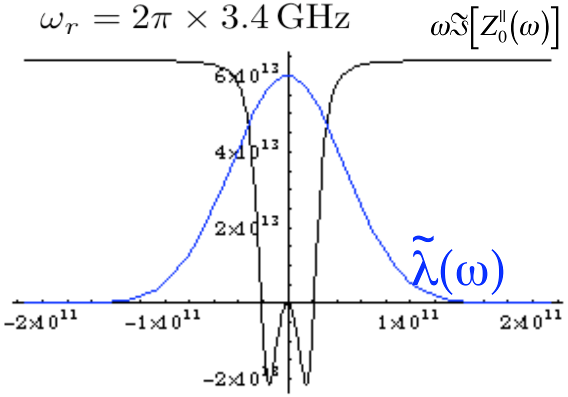
Synchronous phase shift measured with a streak camera in the Australian Synchrotron.

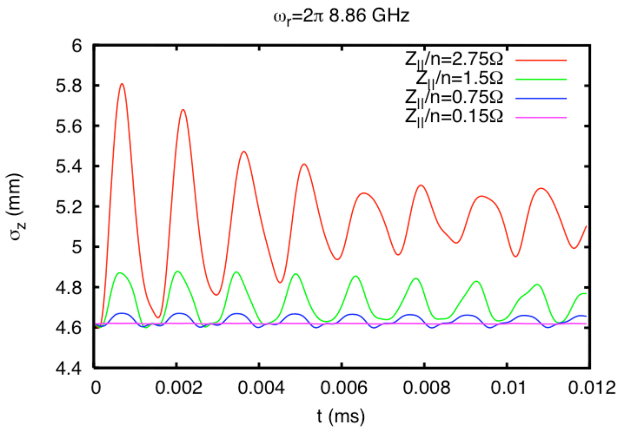
Example of estimations: **Effects of the longitudinal impedance in ALBA**

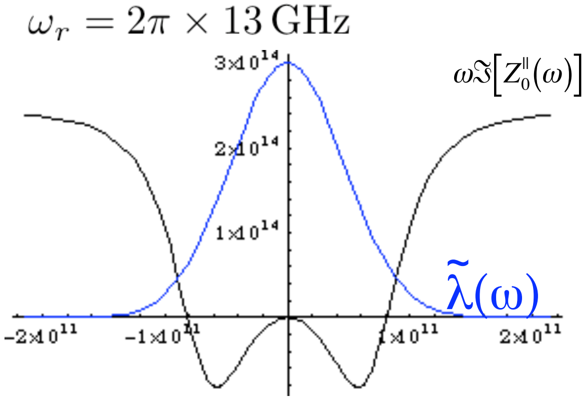
- Simulations can suggest what longitudinal impedance budget is affordable for the future ALBA

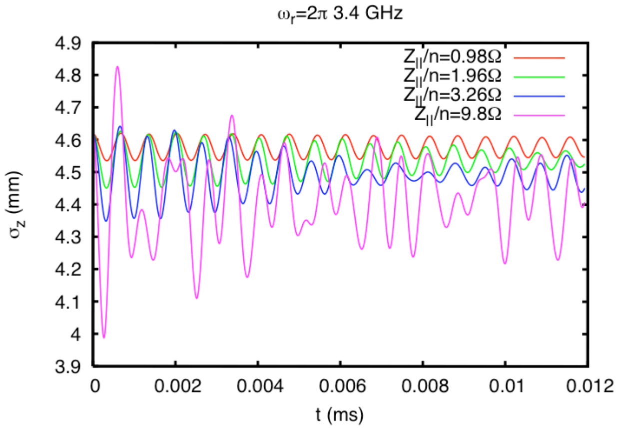
  ⇒ Scan plausible values for R_s. Obviously, the effect of potential well distortion is higher for higher values of R_s

  ⇒ Scan possible values for the frequency. At around 5 GHz the effect of the impedance on the bunch swaps sign (it goes from shortening to lengthening)







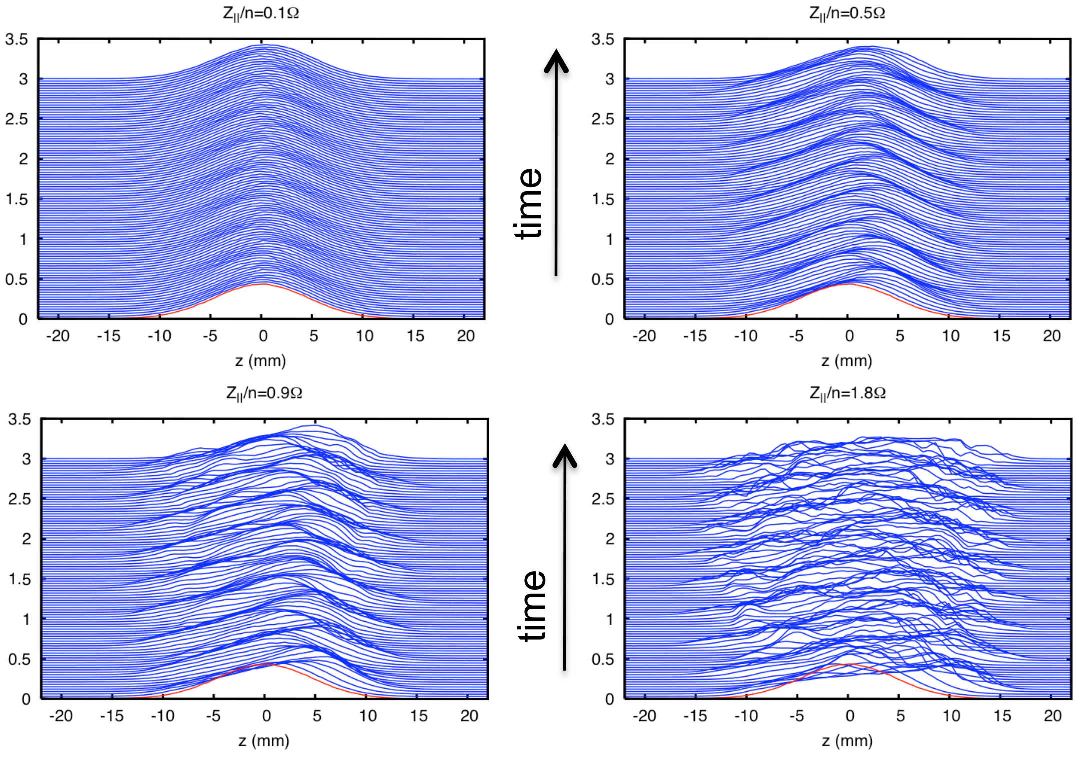


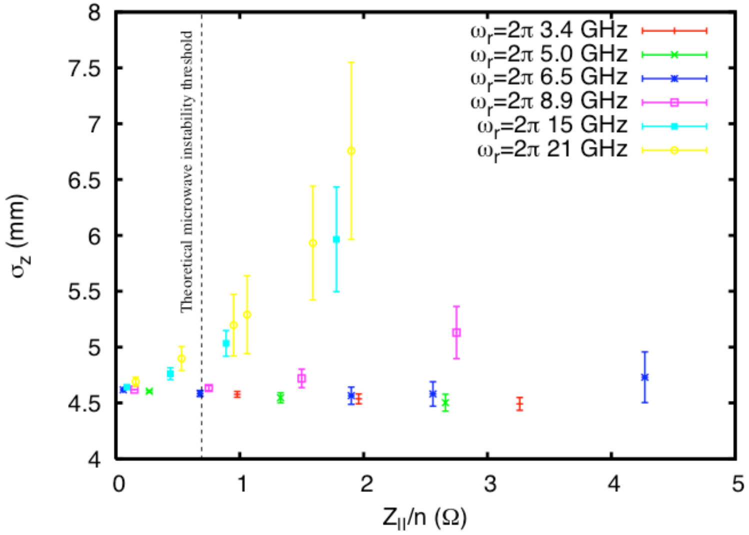
Example of estimations: **Effects of the longitudinal impedance in ALBA**

- Simulations can suggest what longitudinal impedance budget is affordable for the future ALBA

  ⇒ Above a certain value of longitudinal impedance, the microwave instability sets in

  ⇒ We can distinguish the bunch length trend below and above the microwave instability threshold (i.e. potential well and turbulent bunch lengthening regimes) for different frequencies of the BB impedance



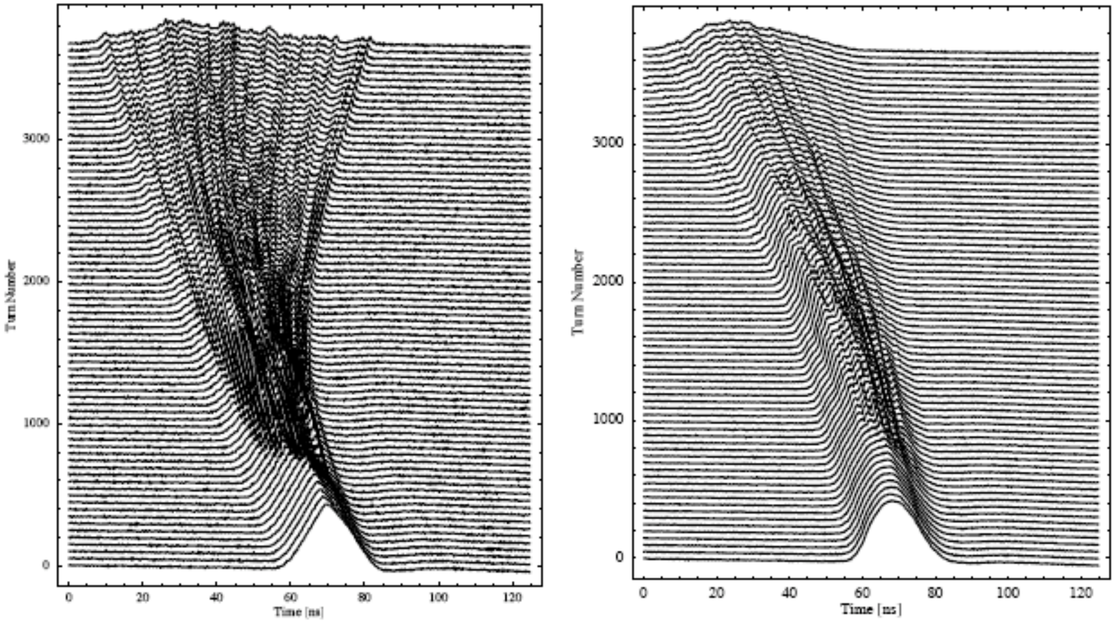


Example of observations: Microwave instability in the SPS

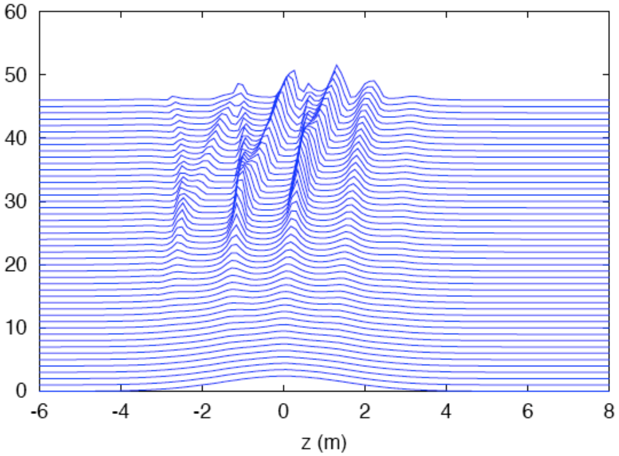
- Microwave instability of a debunching bunch has been used in the SPS to investigate on the spectrum of the longitudinal impedance and try to spot the main frequencies (E. Shaposhnikova, T. Bohl and T. Linnecar)

  ⇒ This allows identifying the main candidates as impedance sources

  ⇒ Long bunch samples better in frequency.



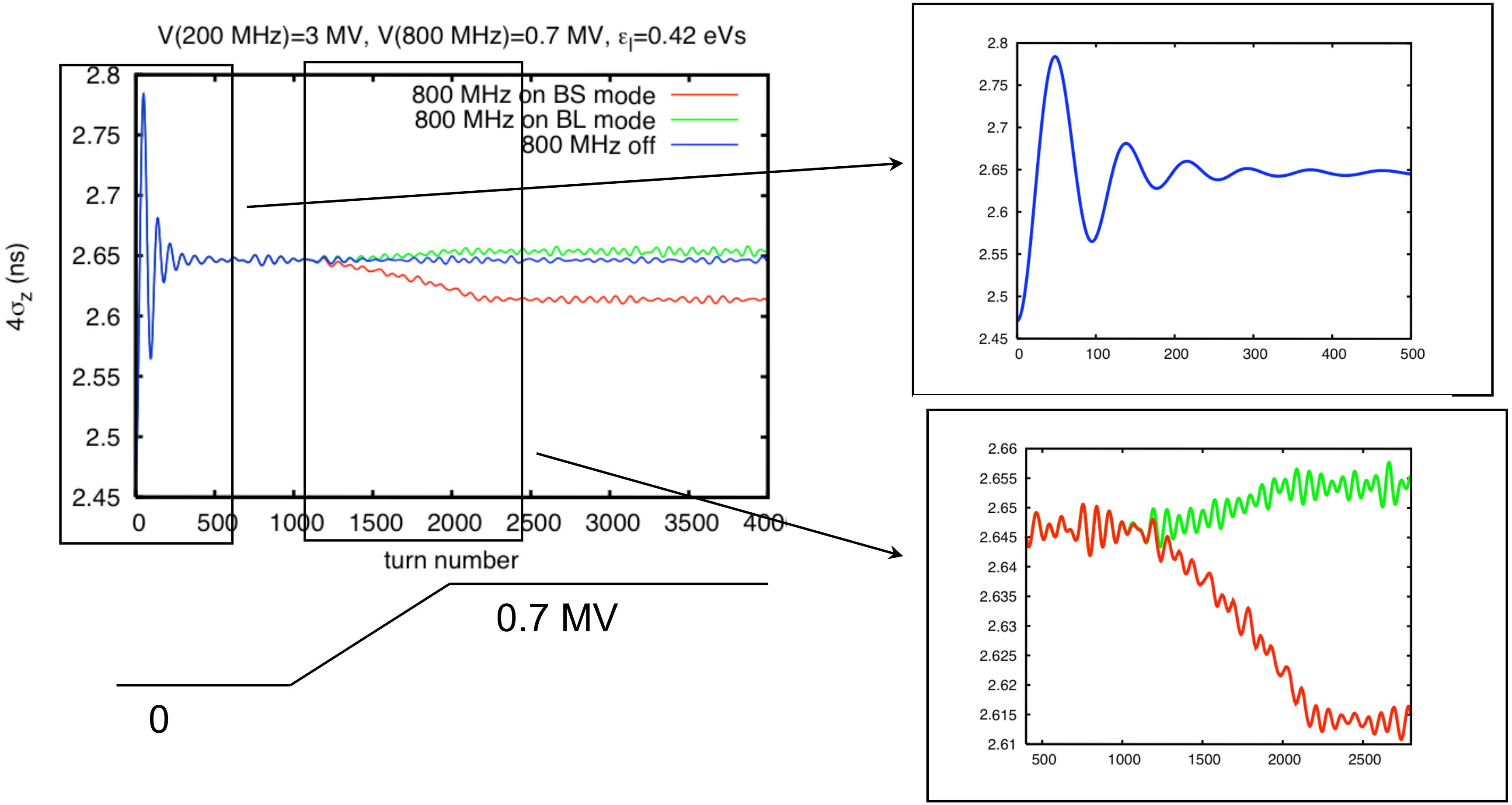
SPS data: below transition energy (left) and above (right)



Simulation with the SPS longitudinal impedance model

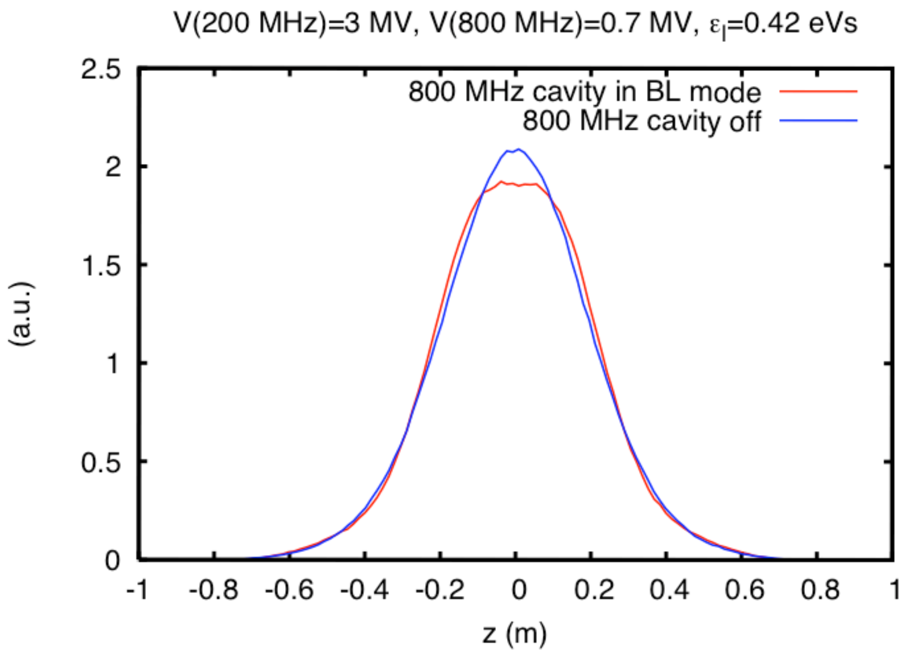
Longitudinal plane: **Production of flat bunches**

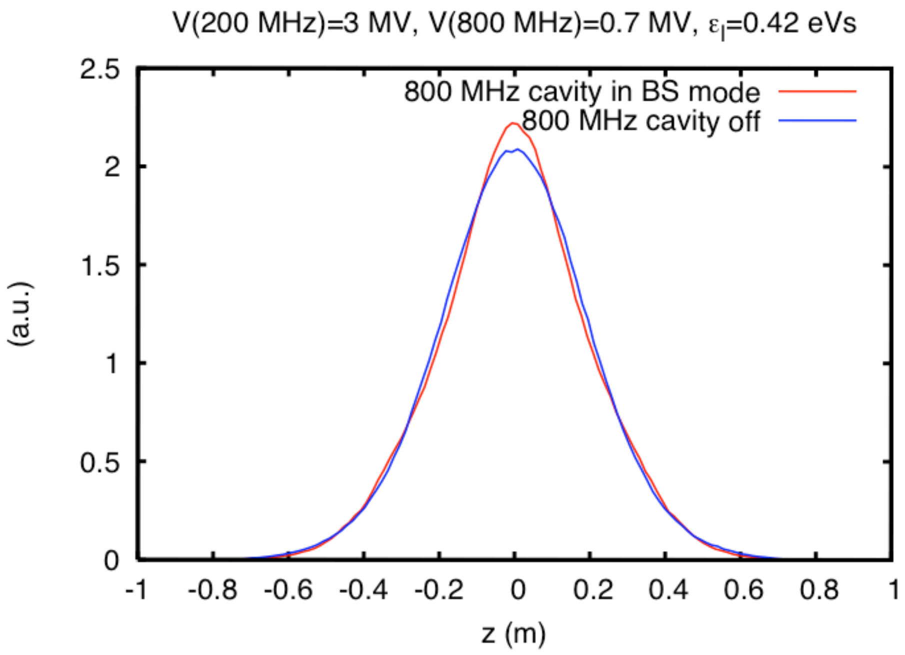
- One option is to use a 2nd harmonic rf system
- The SPS has the main 200 MHz cavities and the 800 MHz, which can be used in Bunch Shortening or Bunch Lengthening modes



Longitudinal plane: **Production of flat bunches**

- The bunch shape when both rf systems are active clearly shows the expected trends





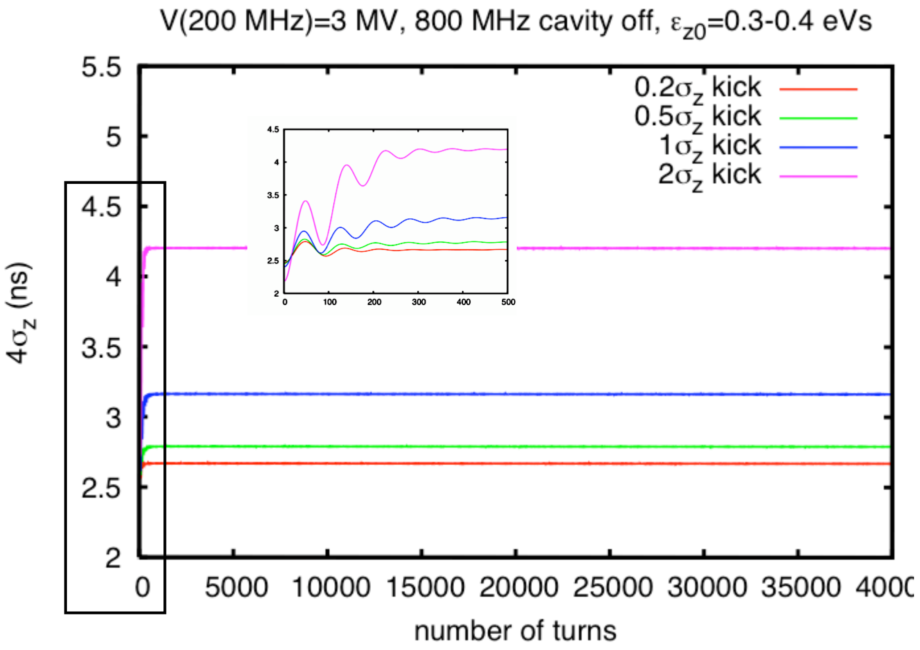
Importance of this option:

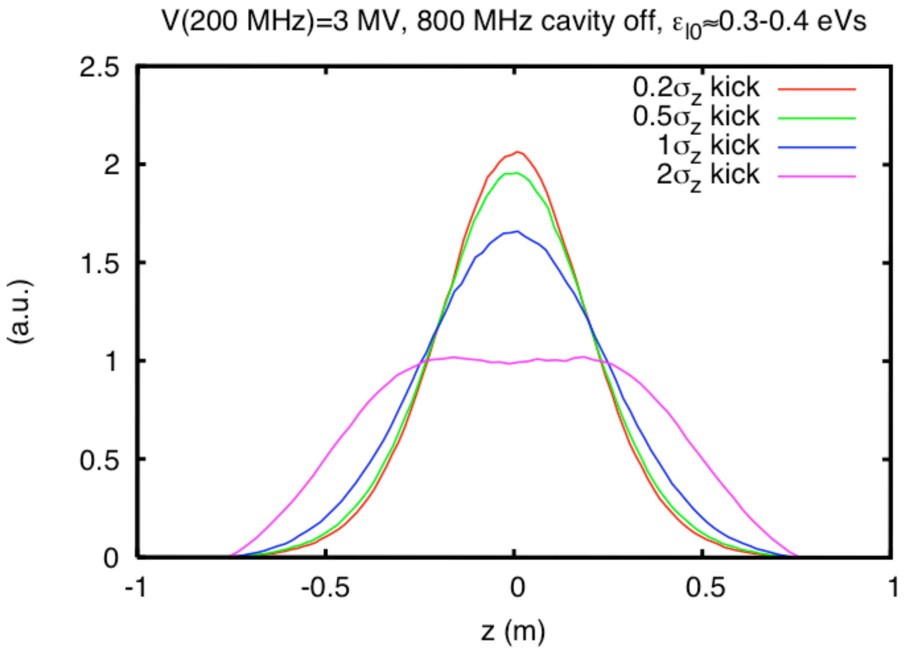
→ SPS: The 800 MHz cavity is used in BS mode in normal operation to keep the beam stable

→ LHC upgrade: Stability studies for a beam in a double rf-system in BL mode (flat bunch)

Longitudinal plane: **Production of flat bunches**

- Another option is to create flat hollow (in longitudinal phase space) bunches
- Experience with hollow bunches shows that problems of stability exist





Importance of this option:

→ LHC upgrade: Simulation studies of stability of flat hollow bunches

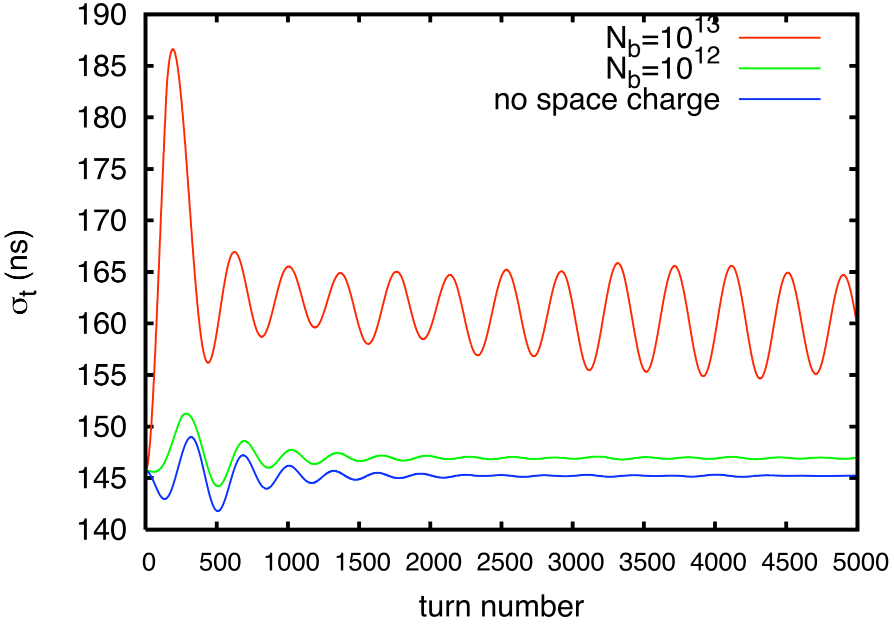
Longitudinal plane: **Production of flat bunches**

- Flattening bunches is also important in low energy for (transverse) space charge reasons.

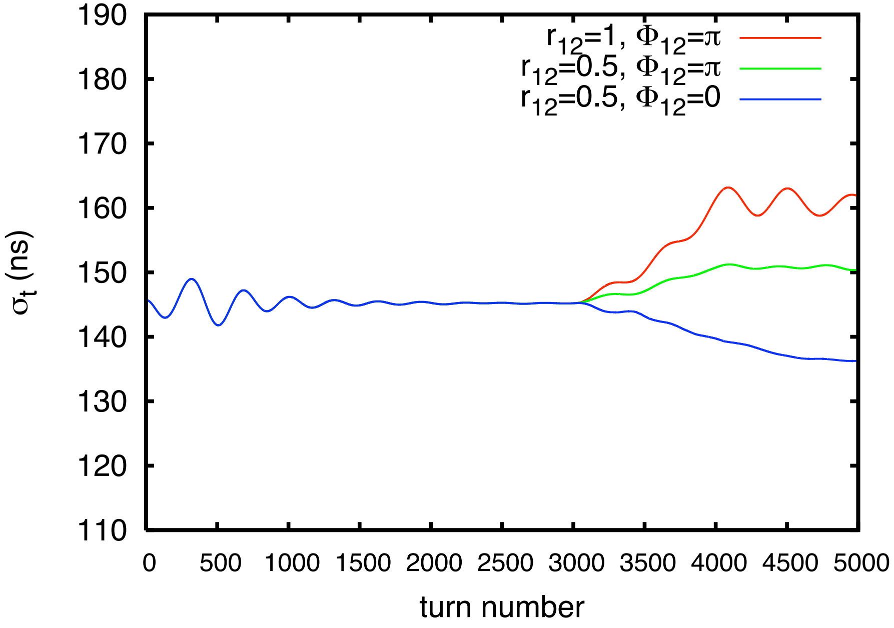
However also longitudinal space charge can then play a role

- For instance, the PSB uses the two rf systems with h=1 and h=2 to accelerate higher intensities

Single rf (h=1): effect of space charge
⇒ space charge lengthens the bunch at equilibrium because we are below transition
⇒ space charge causes the oscillations of mismatching to be undamped



Switch on h=2 system at n=3000
⇒ in BL mode, with two possible ratios between voltages
⇒ in BS mode, with a ratio 0.5 between the two voltages

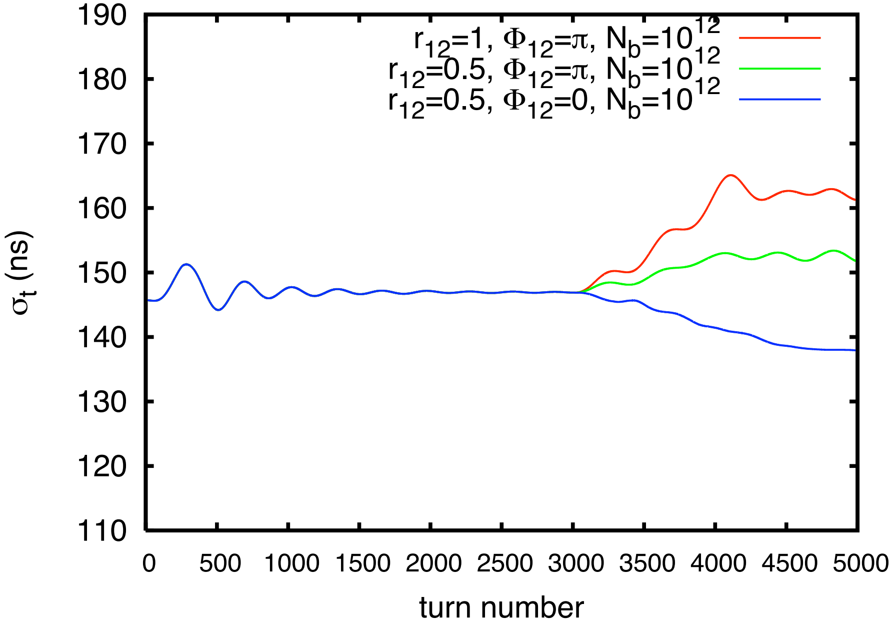


Longitudinal plane: **Production of flat bunches**

- Combined study 2nd harmonic + space charge

Double rf (h=1 + h=2): weak space charge

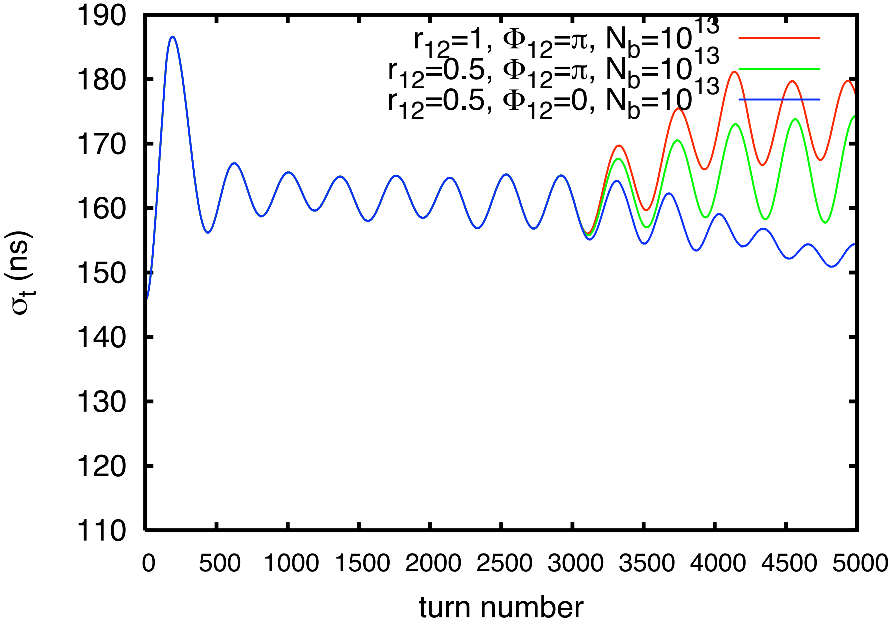
⇒ The final result does not seem to differ much from the case in which space charge was neglected



Double rf (h=1 + h=2): strong space charge

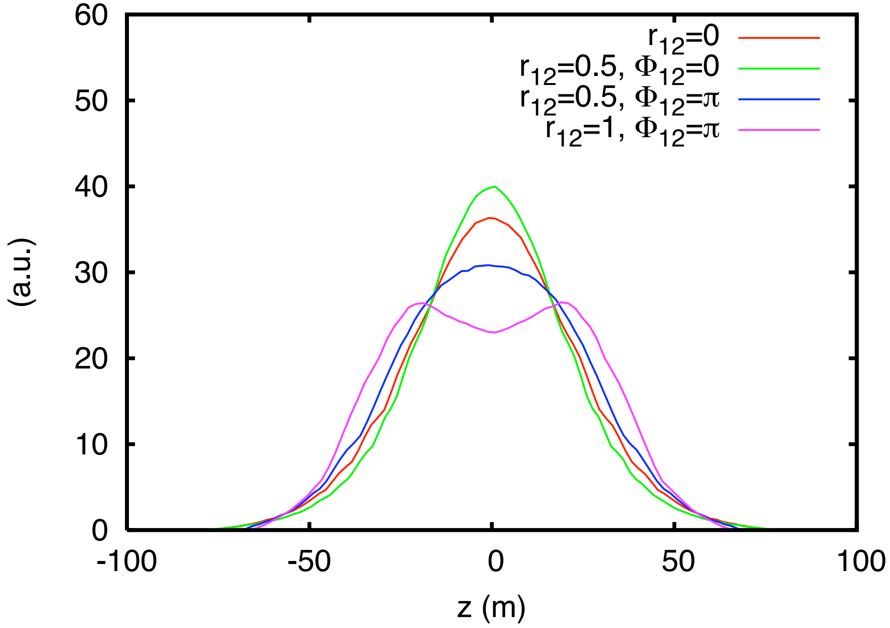
⇒ for higher bunch current the effect of space charge is evident

⇒ oscillations from mismatching are undamped. Feedback (phase loop) needed



Double rf (h=1 + h=2): no long. space charge

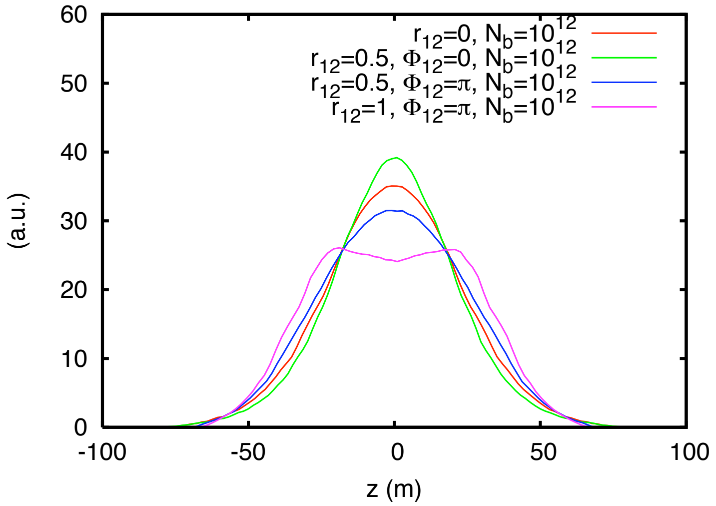
⇒ The effect of the second harmonic on the peak line density (and consequently the beneficial effect against transverse space charge) is evident

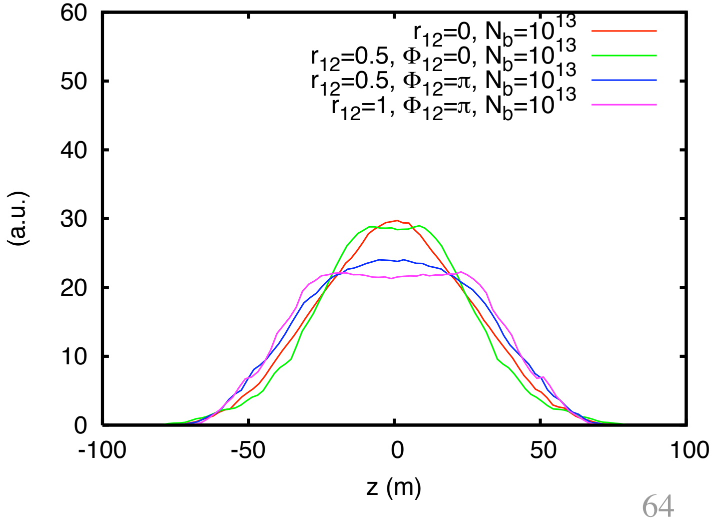


Double rf (h=1 + h=2): weak space charge

⇒ The effect is still significant in not too strong space charge regime, even if the peaks tend to be spread out

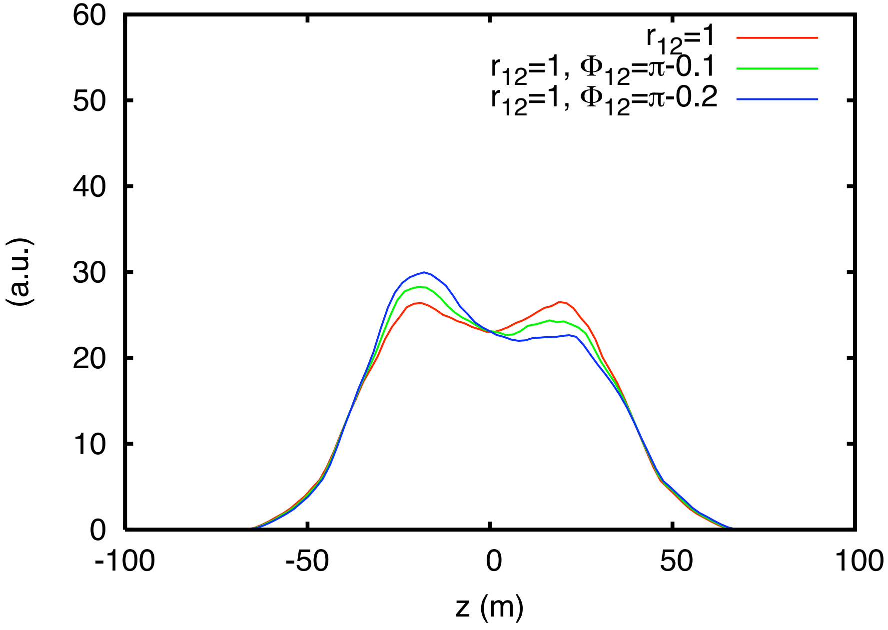
⇒ In high intensity longitudinal space charge causes a loss in peak reduction efficiency

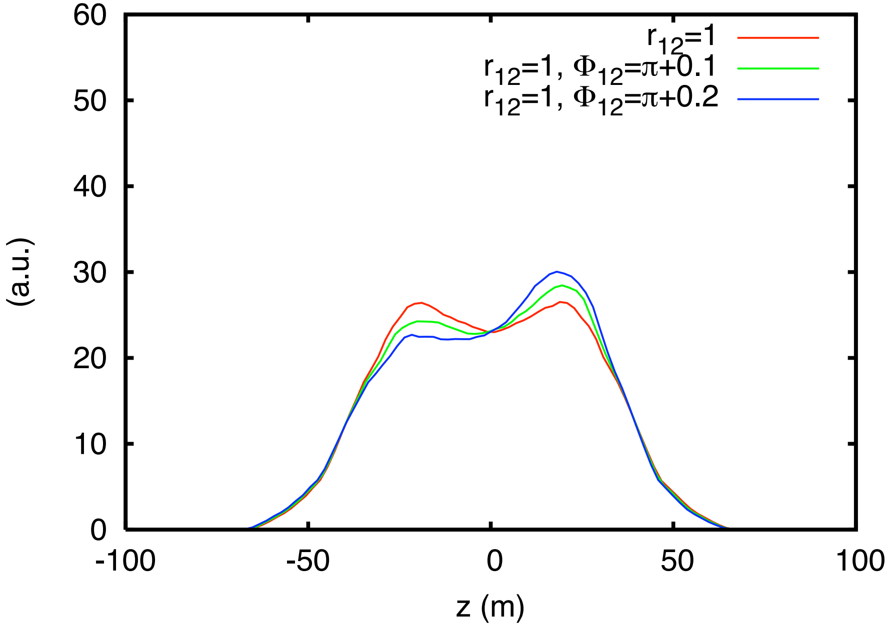




Longitudinal plane: **Production of flat bunches**

- Effect of a phase error in the tuning between the two cavities
- Slight phase errors cause an asymmetric bunch shape, with an enhancement of the right or left peak, depending on the sign of the phase error





Measurements or estimations of the impedance of a machine: **Summary**

- **Transverse:**

⇒ Use growth rates of the mode l=0 of the head-tail instability to estimate the real part of the impedance

→ scan in chromaticity allows for a frequency scan of the impedance spectrum

⇒ Use onset of TMCI and bunch evolution under the effect of a TMCI

⇒ Use coherent tune shift to measure the low frequency imaginary part of the impedance

- **Longitudinal:**

⇒ Several ways to determine the low frequency imaginary part

→ measure the incoherent quadrupole frequency shift for synchrotron oscillations

→ measure bunch lengthening or momentum spread widening

⇒ Real part related to

→ energy loss, which can be estimated by measuring the synchronous phase shift

→ onset of microwave instability.

✓ The rise time relates to the magnitude of the impedance

✓The frequencies involved in the measured evolution also help find possible candidates for main sources of impedance