

# 1- and 2-wire transverse impedance measurement or simulation

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See <http://www.slac.stanford.edu/cgi-wrap/getdoc/slac-ap-110.pdf> (Heifets et al.)  
and  
<http://cdsweb.cern.ch/record/702715/files/sl-note-2002-034.pdf> (Tsutsui)

## ◆ For more precise analyses of machine impedances and related collective effects, several quantities are needed

- $Z_{x,dip}$  (dipolar or driving)
- $Z_{x,quad}$  (“quadrupolar” or detuning)
- $Z_{y,dip}$  (driving)
- $Z_{y,quad}$  (detuning)

$Z_{quad}$  is the same for both transverse planes  
⇒ BUT, 2 differences for its use: 1)  
Subtracts in x and adds in y; 2) Impedances  
are weighted by the betatron function  
in each plane

## “General” definition of transverse impedances (1/2)

- ◆ **Axi-symmetric structures**  $\Rightarrow$  A current density with some azimuthal Fourier component creates electromagnetic fields with the same azimuthal Fourier component

$$\bar{Z}_m = -\frac{1}{Q^2} \int dV \bar{E}_m \cdot \bar{J}_m^*$$

Usual definition of the longitudinal impedance ( $m=0,1,2,\dots$ )

- ◆ **Non axi-symmetric structures**  $\Rightarrow$  A current density with some azimuthal Fourier component may create an electromagnetic field with various different azimuthal Fourier components  $\Rightarrow$  A more general beam coupling impedance is defined in order to treat coupling of different azimuthal Fourier components

$$Z_{m,n} = -\frac{1}{Q^2} \int dV E_m \cdot J_n^*$$

More “general” definition of the longitudinal impedance ( $m,n = 0, \pm 1, \pm 2,\dots$ )

## “General” definition of transverse impedances (2/2)

- ◆ The “general” transverse impedances  $Z_{x,y}$  (in  $\Omega!$ , i.e. not normalized by the transverse displacement) on a test particle at  $(x_2 = a_2 \cos\theta_2, y_2 = a_2 \sin\theta_2)$  from a source at  $(x_1 = a_1 \cos\theta_1, y_1 = a_1 \sin\theta_1)$ , are given by (to 1<sup>st</sup> order)

$$k Z_x = \left( Z_{0,1} + Z_{0,-1} \right) + x_1 \bar{Z}_x + j y_1 \left( -Z_{1,-1} - Z_{1,1} + Z_{-1,-1} + Z_{-1,1} \right) + 2 \left( Z_{0,2} + Z_{0,-2} \right) x_2 + 2 \left( Z_{0,2} - Z_{0,-2} \right) j y_2$$

$$k Z_y = j \left( Z_{0,1} - Z_{0,-1} \right) + y_1 \bar{Z}_y + j x_1 \left( -Z_{1,-1} + Z_{1,1} - Z_{-1,-1} + Z_{-1,1} \right) - 2 \left( Z_{0,2} + Z_{0,-2} \right) y_2 + 2 \left( Z_{0,2} - Z_{0,-2} \right) j x_2$$

with

$$k = \omega / c$$

$$\bar{Z}_x = Z_{1,1} + Z_{1,-1} + Z_{-1,1} + Z_{-1,-1}$$

$$\bar{Z}_y = Z_{1,1} - Z_{1,-1} - Z_{-1,1} + Z_{-1,-1}$$

⇒

$$Z_x^{\text{driving}} = \bar{Z}_x / k$$

$$Z_y^{\text{driving}} = \bar{Z}_y / k$$

$$Z^{\text{detuning}} = -2 \left( Z_{0,2} + Z_{0,-2} \right) / k$$

# 2-wire technique

- ◆ This gives the (usual) dipolar (driving) impedance
- ◆ Method: Measure the longitudinal impedance  $Z$  and deduce the transverse one

$$Z = 2 Z_{ch} \frac{1 - S_{21}}{S_{21}}$$

Characteristic impedance

$$Z_x^{\text{driving}} = \frac{\bar{Z}_x}{k} = \frac{c Z}{\omega (2a)^2}$$

The 2 wires are at  $\pm a$

Scattering parameter

# 1-wire technique (1/4)

- ◆ With 1 wire (at  $x = a \cos\theta$ ,  $y = a \sin\theta$ ), the longitudinal impedance measured (or simulated) is given by (to 2<sup>nd</sup> order)

$$Z = A_1 + a e^{-j\vartheta} A_2 + a e^{j\vartheta} A_3 + a^2 e^{-2j\vartheta} A_4 + a^2 e^{2j\vartheta} A_5 + a^2 A_6$$

with

$$A_1 = Z_{0,0}$$

$$A_2 = Z_{1,0} + Z_{0,-1}$$

$$A_3 = Z_{0,1} + Z_{-1,0}$$

$$A_4 = Z_{2,0} + Z_{1,-1} + Z_{0,-2}$$

$$A_5 = Z_{0,2} + Z_{-1,1} + Z_{-2,0}$$

$$A_6 = Z_{1,1} + Z_{-1,-1}$$

# 1-wire technique (2/4)

- ◆ If there is top/bottom and left/right symmetry (fortunately it is the usual case...), the situation simplifies a lot  $\Rightarrow$  See EPAC06 paper (<http://accelconf.web.cern.ch/AccelConf/e06/PAPERS/THPCH059.PDF>)

- If  $\mathbf{a} = \mathbf{x}_0$  and  $\theta = 0$

$$\begin{aligned} Z &= A_1 + x_0^2 (A_4 + A_5 + A_6) \\ &= A_1 + x_0^2 \left[ \bar{Z}_x + (Z_{2,0} + Z_{0,2} + Z_{-2,0} + Z_{0,-2}) \right] \end{aligned}$$

Scanning  $x_0$   
gives a parabola

- If  $\mathbf{a} = \mathbf{y}_0$  and  $\theta = \pi/2$

$$\begin{aligned} Z &= A_1 + y_0^2 (-A_4 - A_5 + A_6) \\ &= A_1 + y_0^2 \left[ \bar{Z}_y - (Z_{2,0} + Z_{0,2} + Z_{-2,0} + Z_{0,-2}) \right] \end{aligned}$$

# 1-wire technique (3/4)

⇒ IF  $Z_{2,0} + Z_{-2,0} = Z_{0,2} + Z_{0,-2}$  , then

$$Z_{2,0} + Z_{0,2} + Z_{-2,0} + Z_{0,-2} = 2 \left( Z_{0,2} + Z_{0,-2} \right) = -k Z^{\text{detuning}}$$

Tsutsui showed “approximately” that for a 2D thick metallic boundary case:  $Z_{-m,-n} = Z_{n,m} \Rightarrow$  In this case, the above formula is indeed valid

**Remark 1:** Only the “generalized impedances” can be obtained but not the driving and detuning impedance separately! This cannot be used like this in codes like HEADTAIL for instance

**Remark 2:** The dipolar (driving) impedances need to be obtained from the 2-wire technique

# 1-wire technique (4/4)

- ◆ If there is NO top/bottom or left/right symmetry, the situation is more involved:

- By scanning  $a$  and  $\theta$  (i.e. measuring  $Z$  for different values of  $a$  and  $\theta$ ),  $A_{1,2,3,4,5,6}$  can be found

- Then, using the 2-wire technique the dipolar (driving) impedances can be obtained:

$$Z_x^{\text{driving}} = \bar{Z}_x / k$$

$$Z_y^{\text{driving}} = \bar{Z}_y / k$$

- Then compute  $Z_{1,-1} + Z_{-1,1} = (\bar{Z}_x - \bar{Z}_y) / 2$

- Then, if  $Z_{2,0} + Z_{-2,0} = Z_{0,2} + Z_{0,-2}$ ,

$$Z^{\text{detuning}} = \frac{Z_x^{\text{driving}} - Z_y^{\text{driving}}}{2} - \frac{A_4 + A_5}{k}$$



# Conclusion (and reminder)

- ◆ Both 1-wire and 2-wire techniques are required (in asymmetric structures) to obtain all the information needed to correctly understand/describe the collective effects in accelerators
- ◆ With 2 wires the transverse dipolar (driving) impedances are obtained
- ◆ With 1 wire (scanning the wire position), and using the driving impedances measured with 2 wires, the detuning impedance can be deduced (IF a certain condition is fulfilled  $\Rightarrow$  Still to be checked in which cases this relation is satisfied or not)
- ◆ What about the coupling terms (see page 3)??? Can they be important in some cases???